

# Finite temperature QCD with chiral (overlap) fermions

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in collaboration with **Sz. Borsanyi, T. Kovacs,  
K. Szabo, Z. Fodor**

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PoS LATTICE2023 (2024) 179, arXiv: 2401.02750

+ new results



**New Trends in Thermal Phases of QCD,  
Bratislava, 2026**

# Outline

- Lattice setup
- Topological susceptibility  $\chi$
- Nature of the chiral phase transition
- Spectrum of the Dirac operator and localisation

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  - Two boson fields with mass  $m_B a = 0.54$
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  - Everything is preliminary!
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# Topological susceptibility from simulations at fixed $Q$

## Correlator method

$$\lim_{|z| \rightarrow \text{large}} \langle q(z)q(0) \rangle = \frac{1}{V} (Q^2/V - \chi) + O(V^{-2}) + O(e^{-m_\eta' z})$$

$q(z)$ : topological charge density

[Aoki, Fukaya, Hashimoto, Onogi, 2007]

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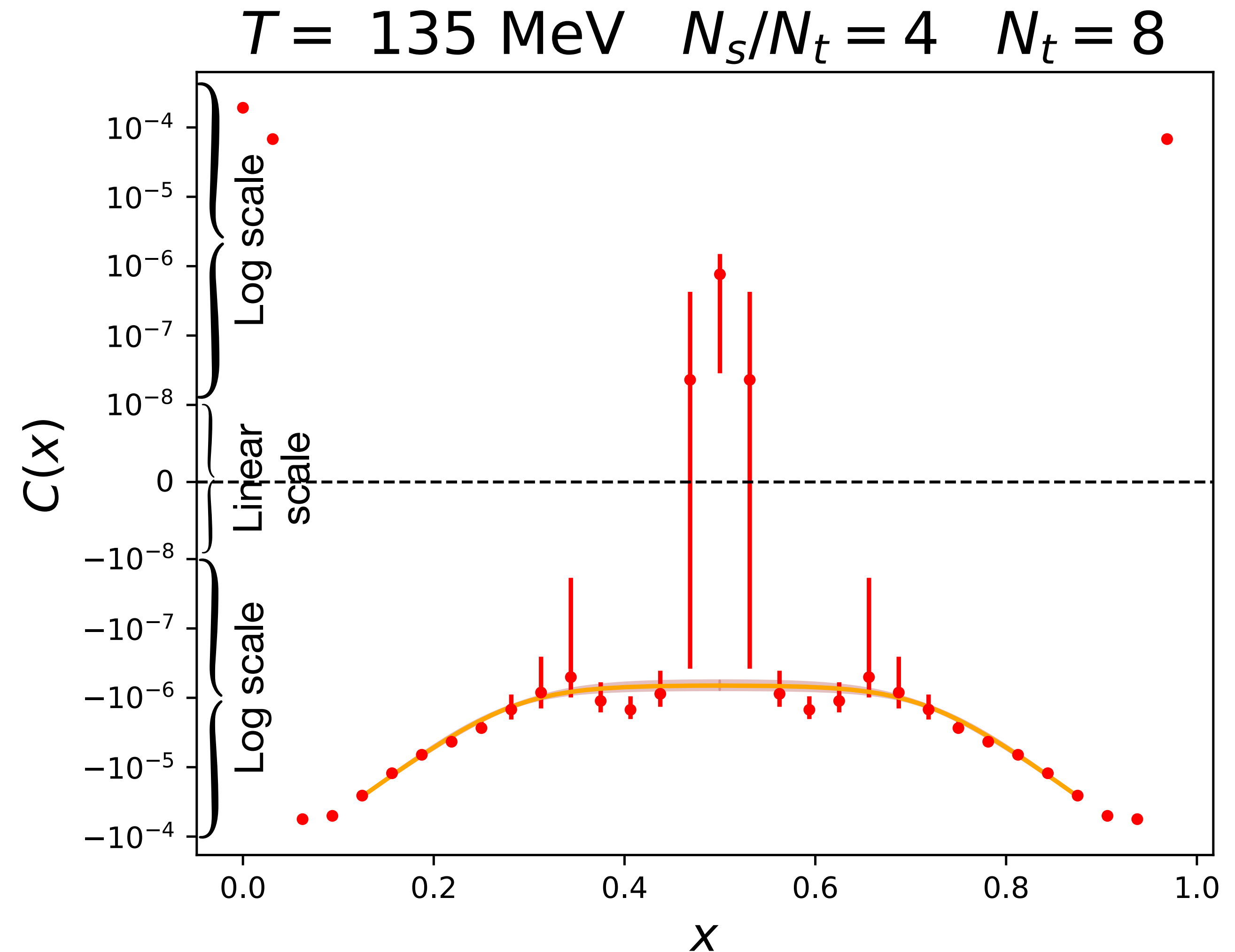
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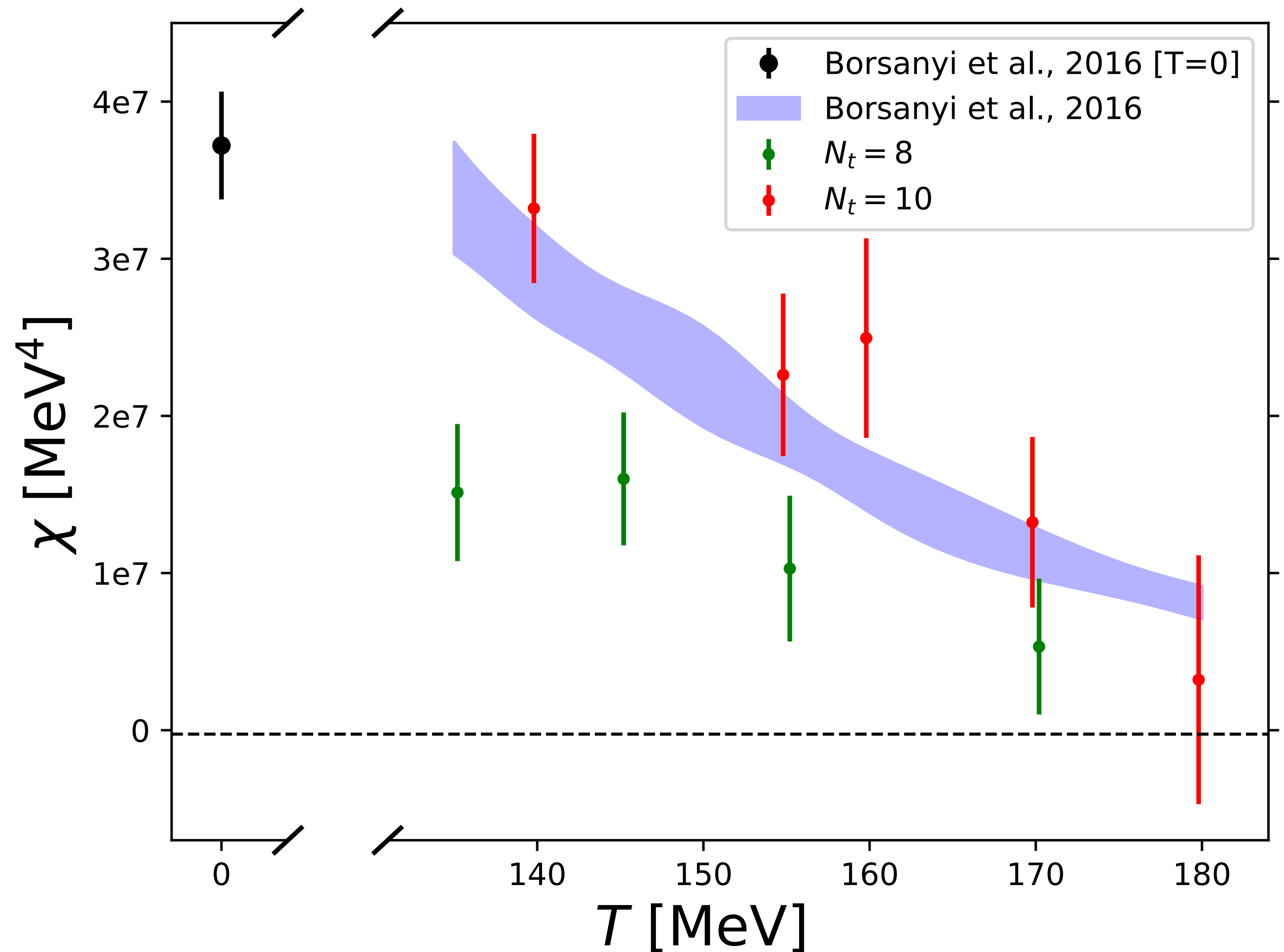
- More stable for  $N_s/N_t \gtrsim 4$



# Topological susceptibility from simulations at fixed $Q$

## Correlator method

- Noisy
- Consistent with [\[Borsanyi et al., 2016\]](#)
- Local topological fluctuations
- $N_t = 8$ : lattice artefacts?



# Chiral observables

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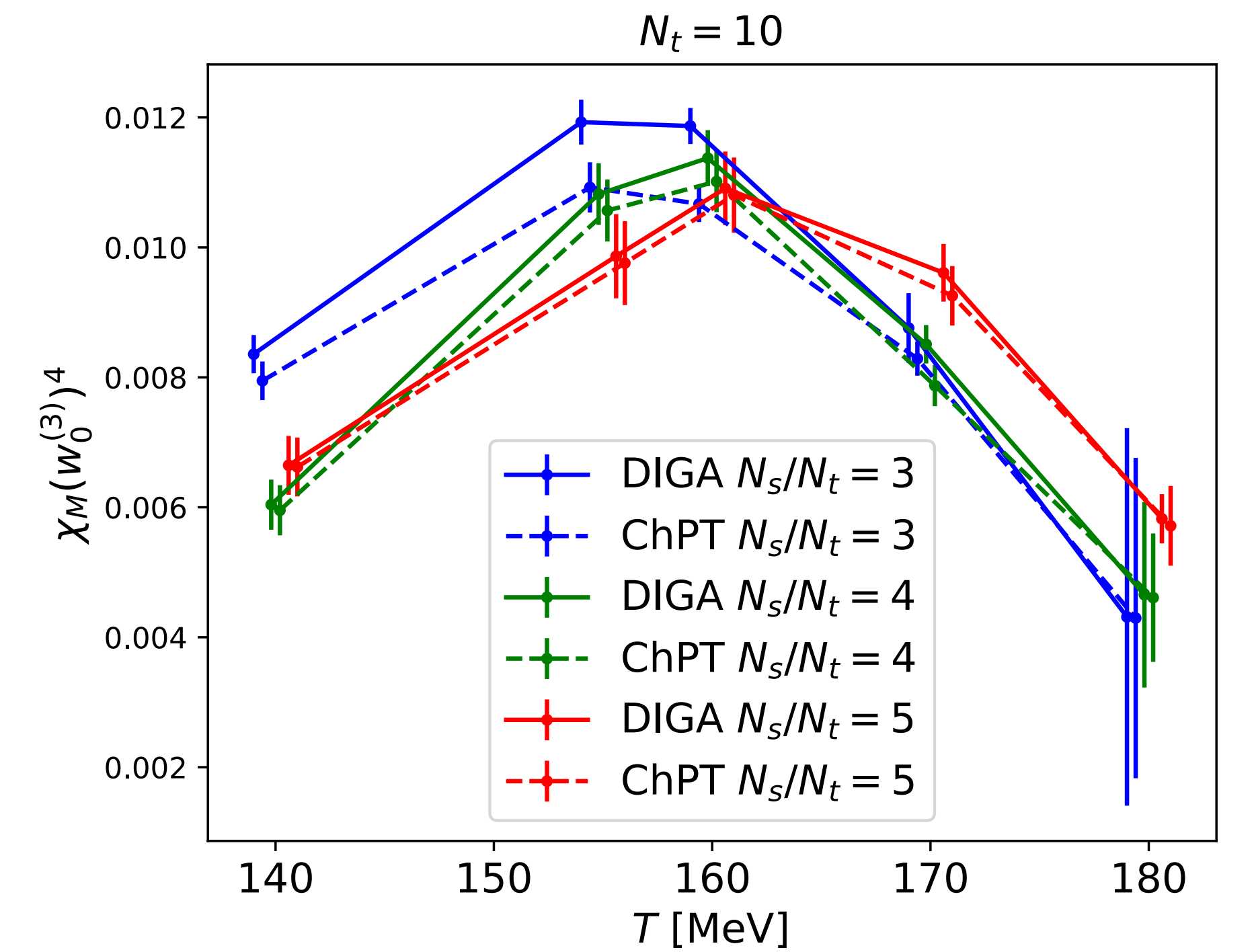
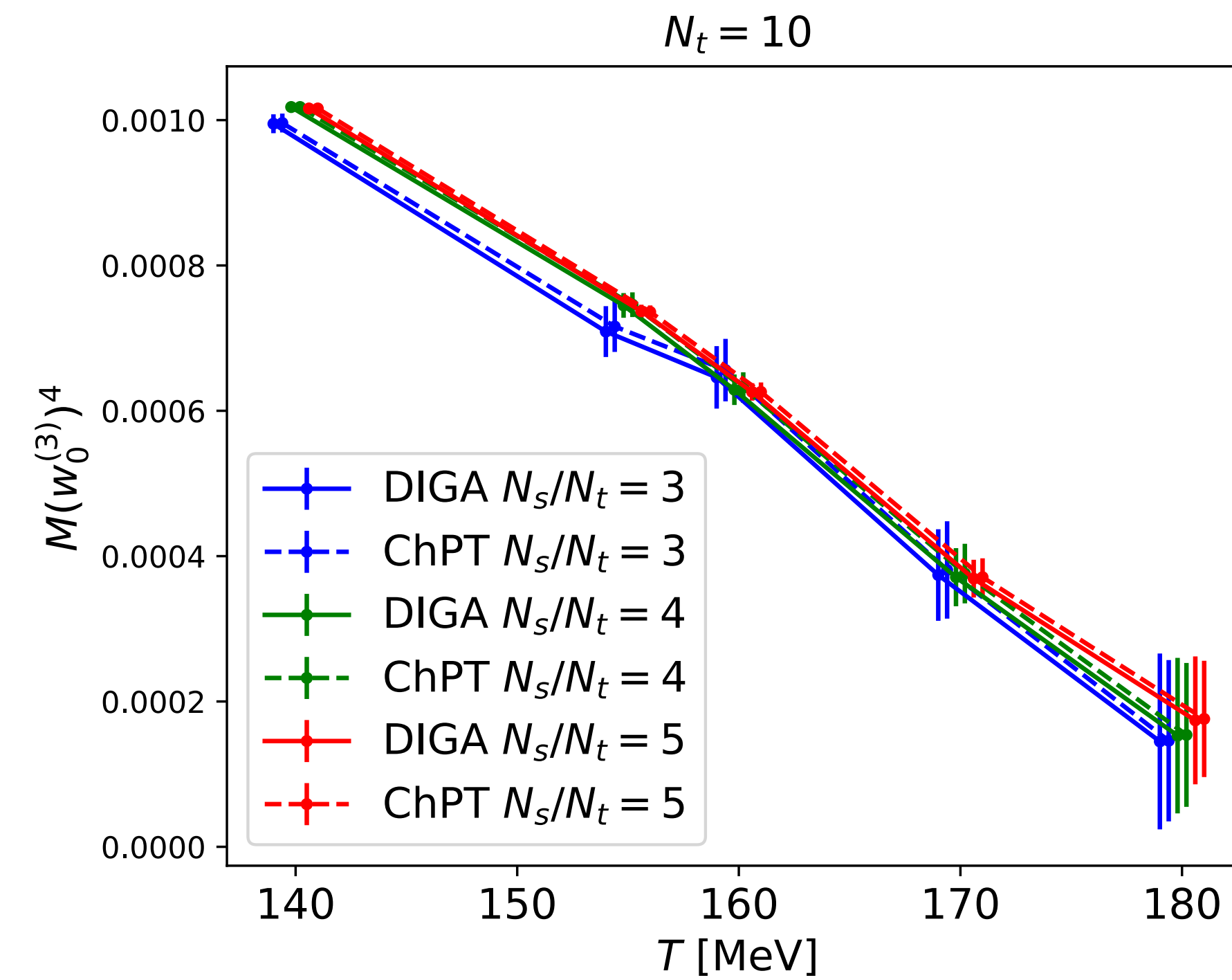
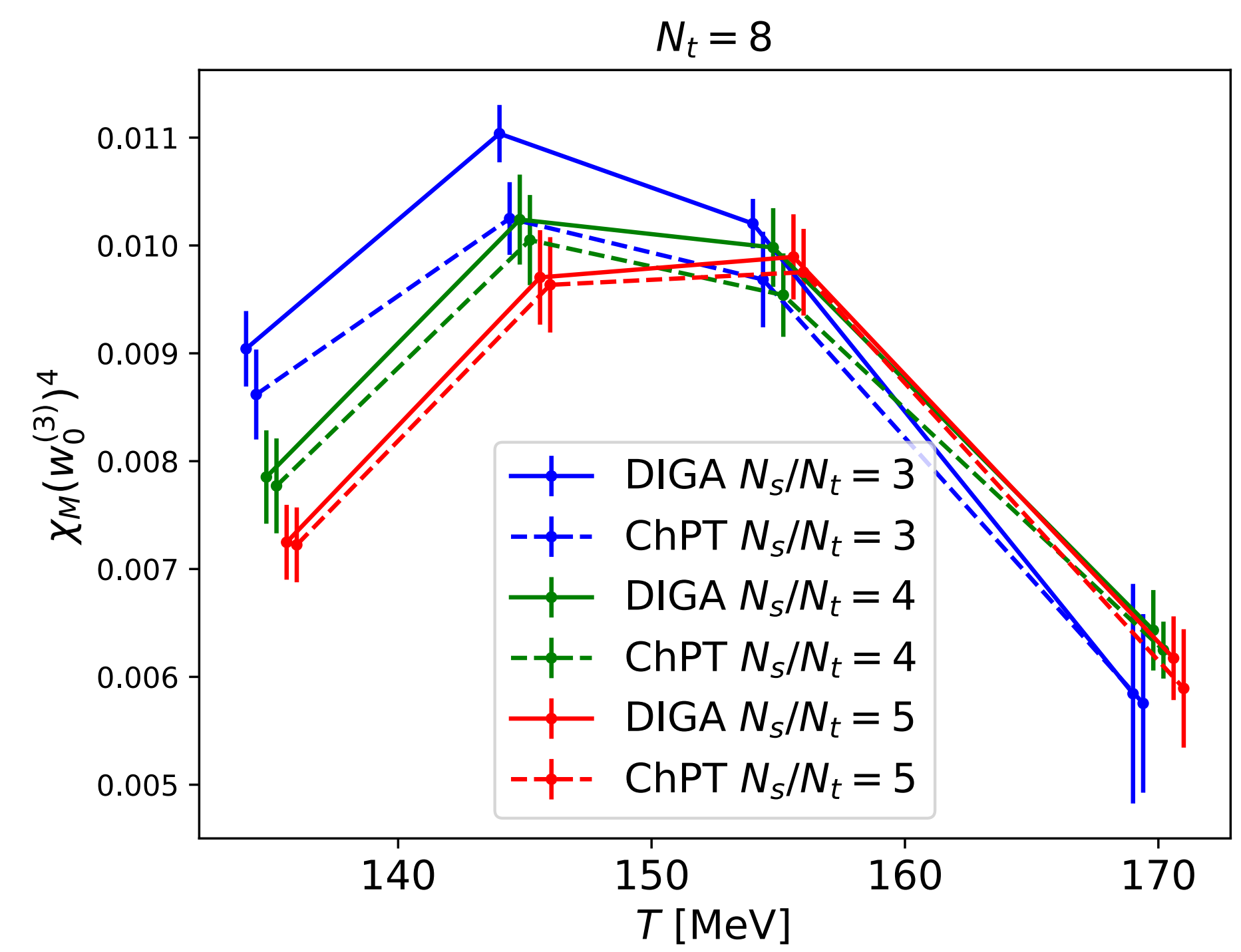
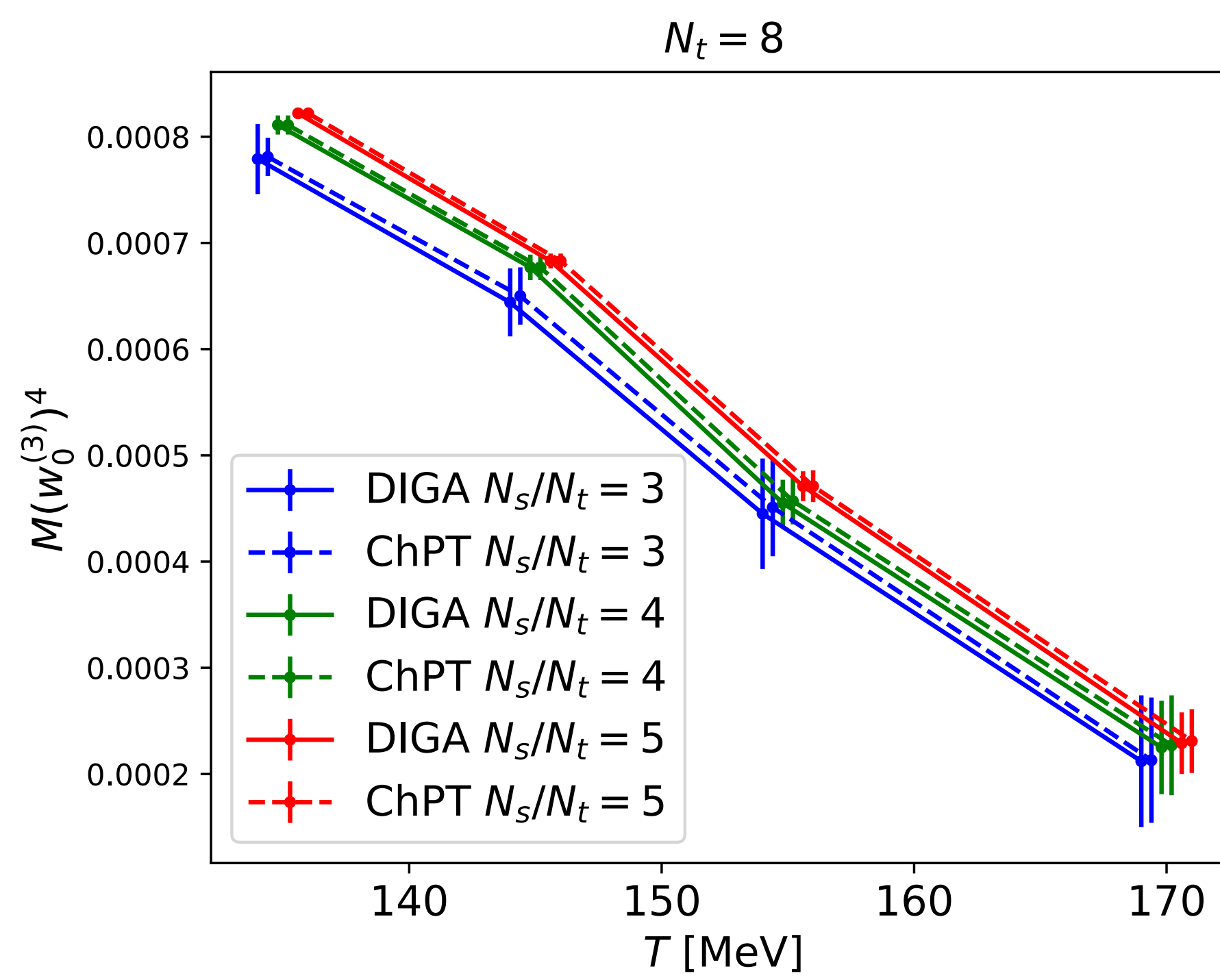
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- Weights of different sectors:
  - LO ChPT:  $Z_Q \sim \int d\theta \frac{I_1(a \cos(\theta/2))}{a \cos(\theta/2)} e^{iQ\theta}$  [Leutwyler, Smilga, 1992]
  - DIGA:  $Z_Q \sim I_Q(\chi V)$

# Chiral phase transition

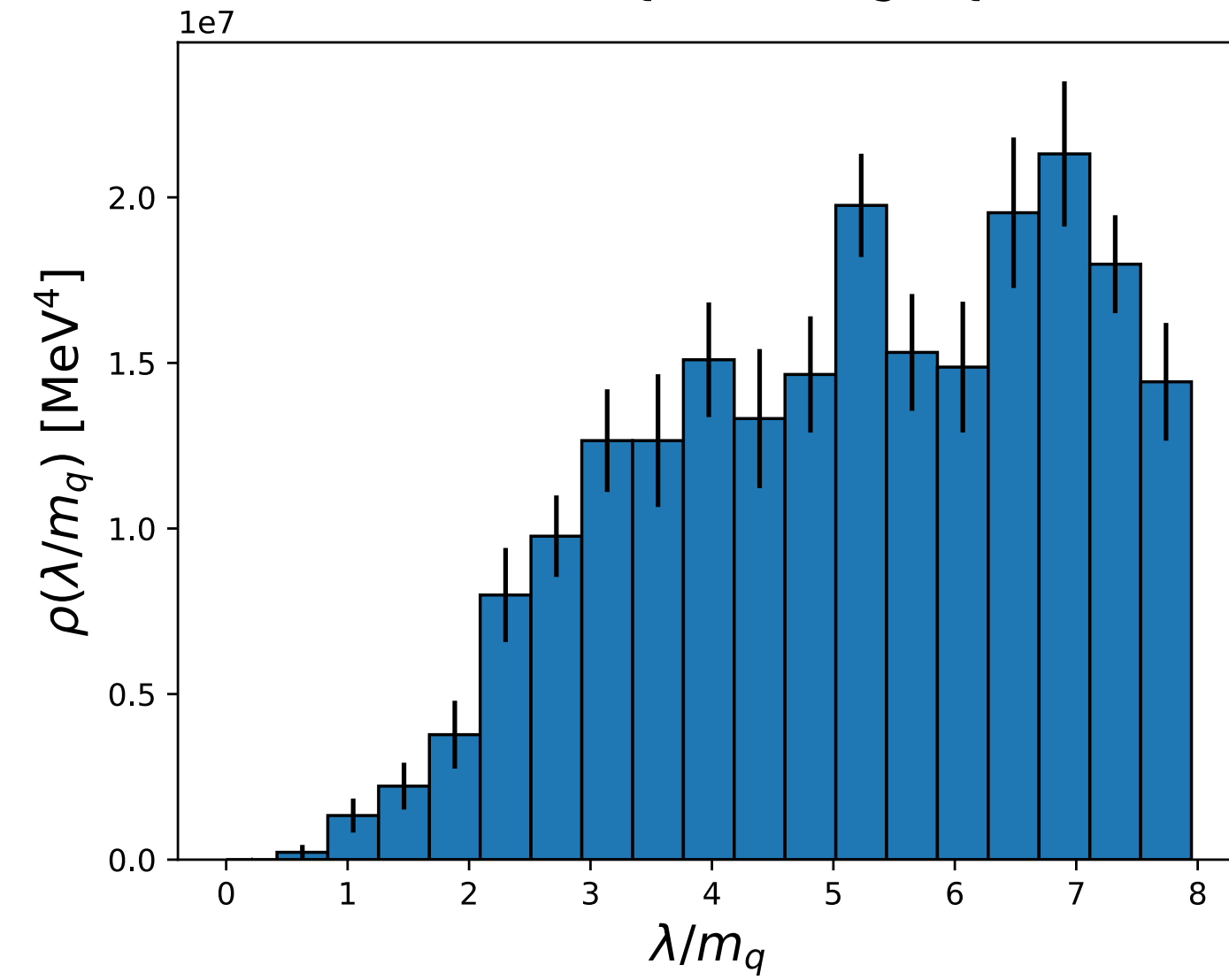


# Dirac operator spectrum, $T = 145$ MeV

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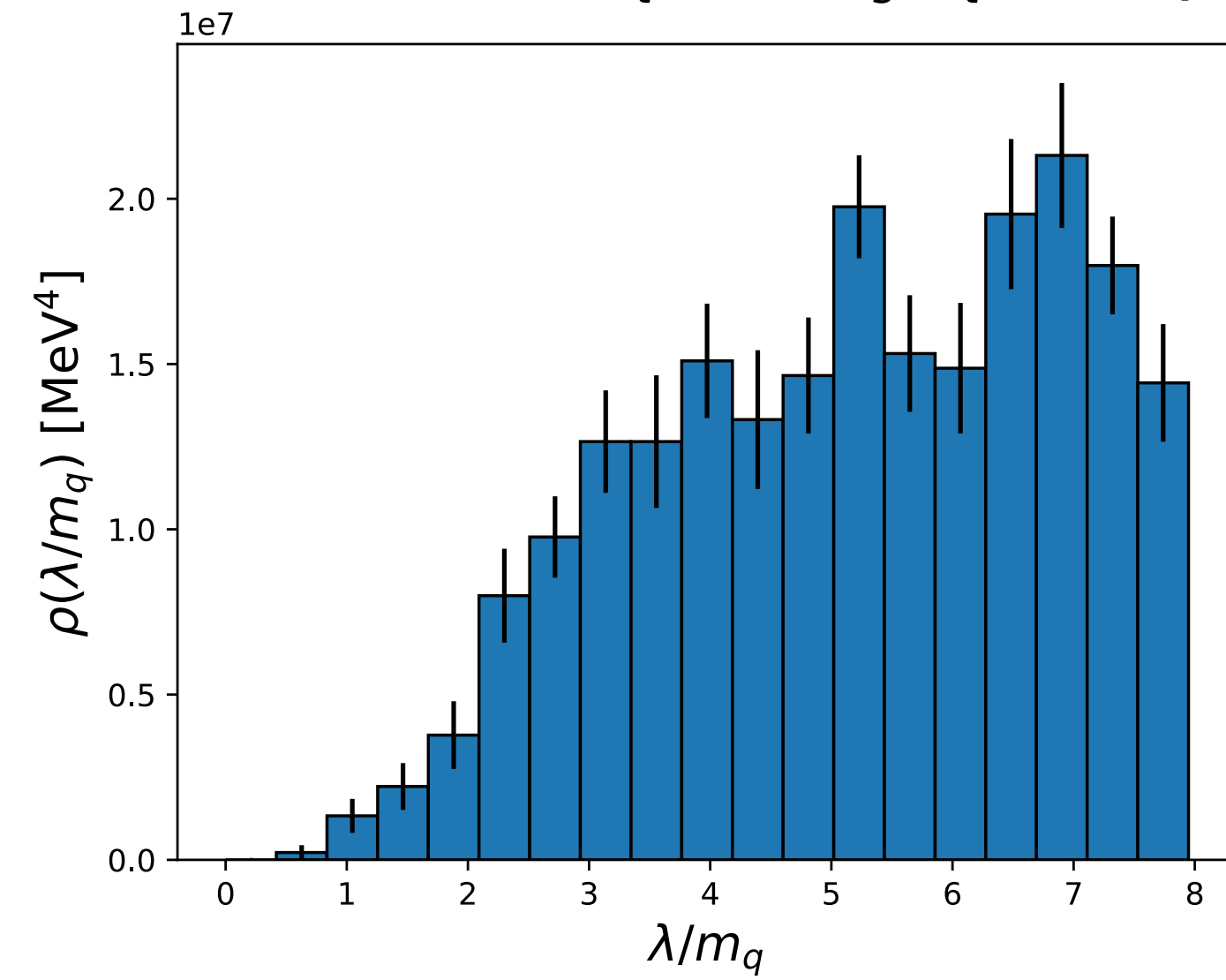
$T = 145.0 \text{ MeV}$   $N_t = 8$   $N_s/N_t = 2$   $Q = 0$



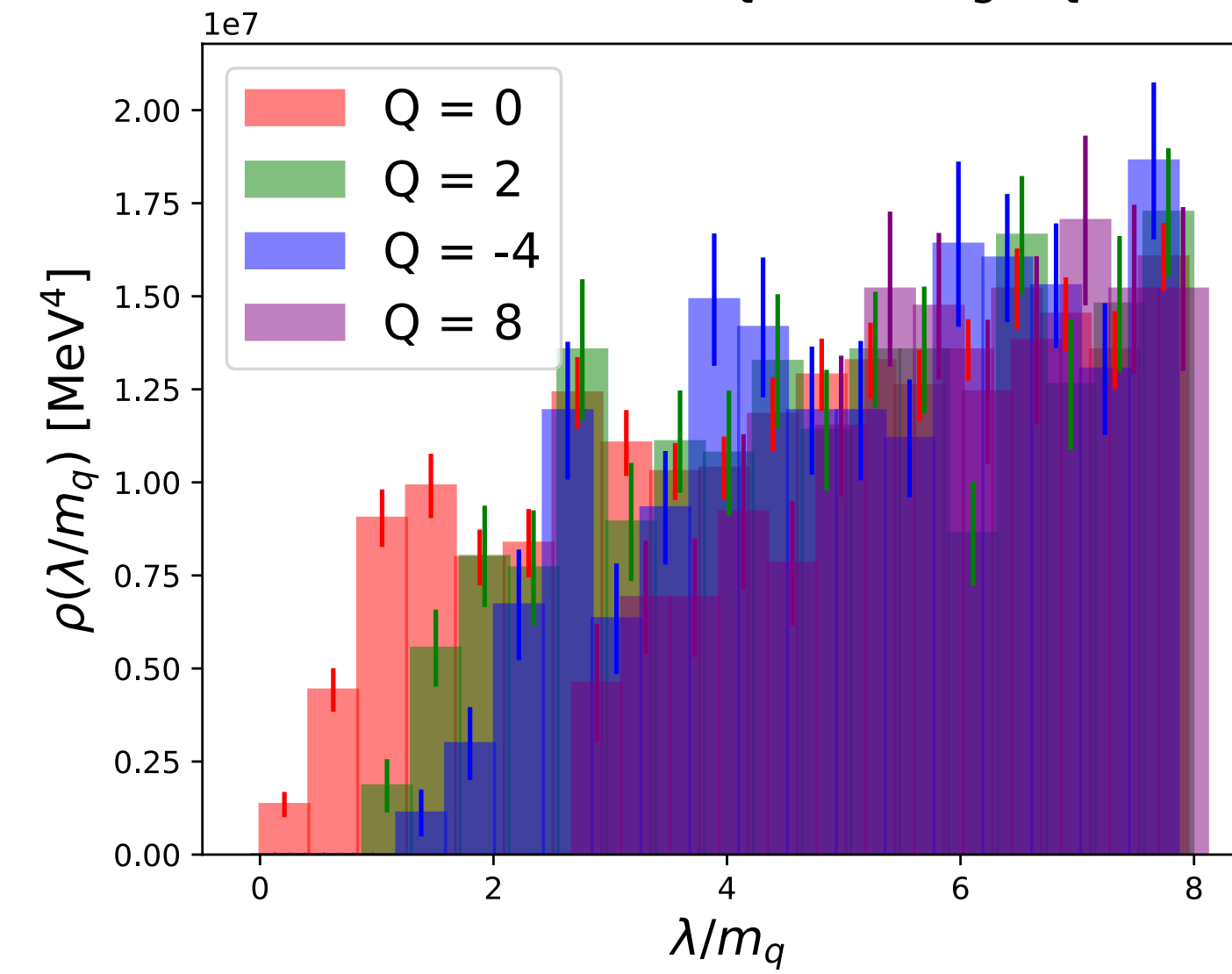
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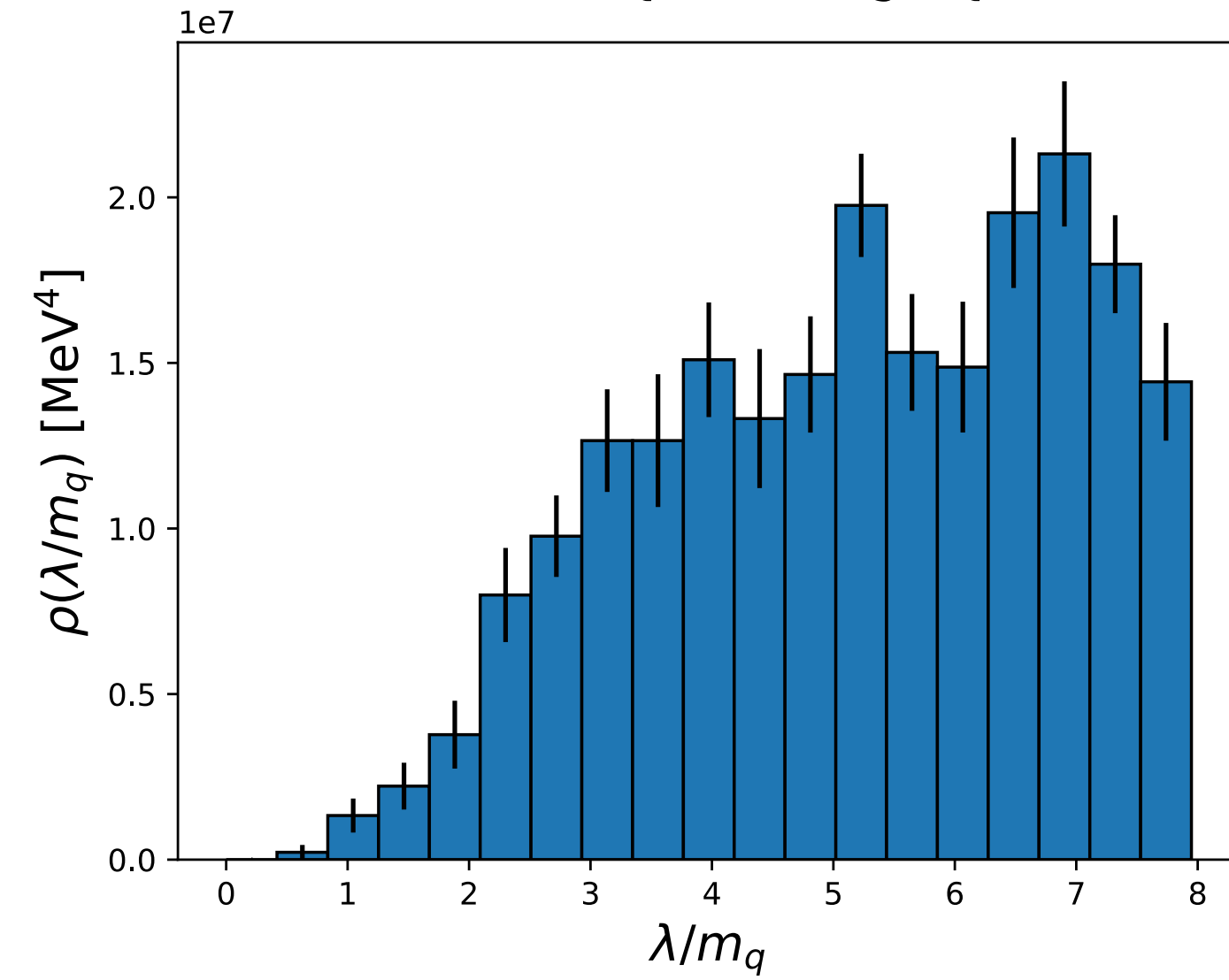


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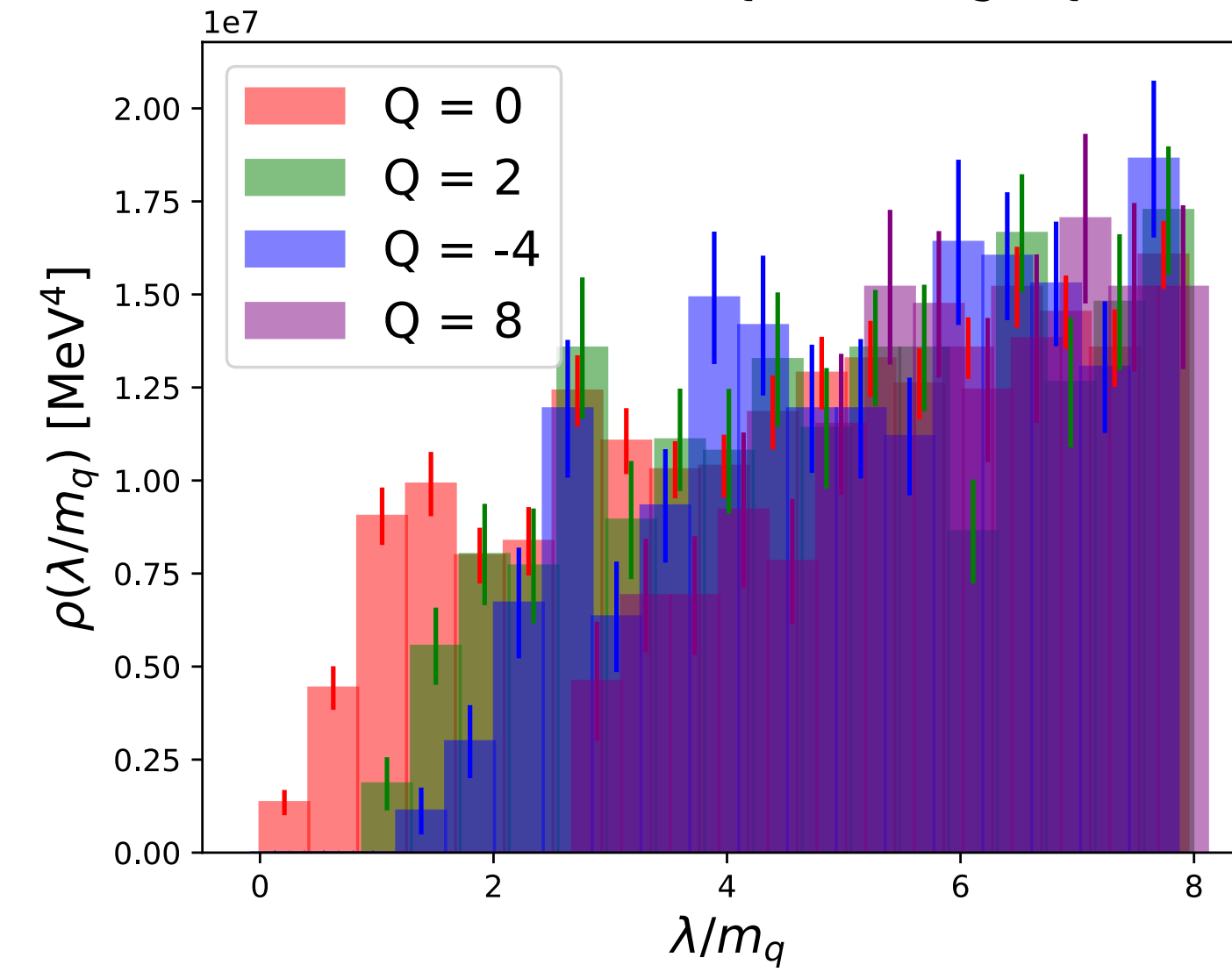
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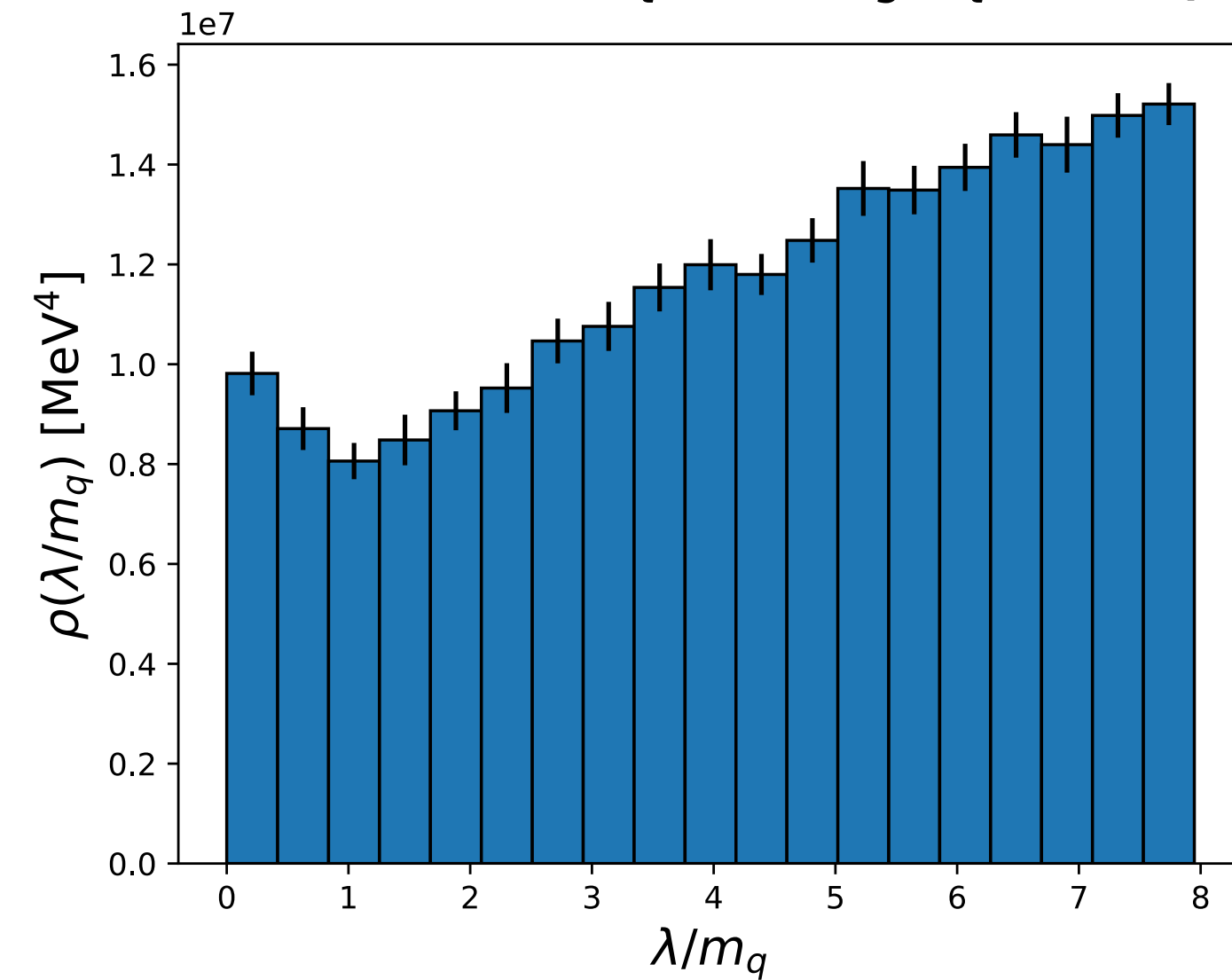
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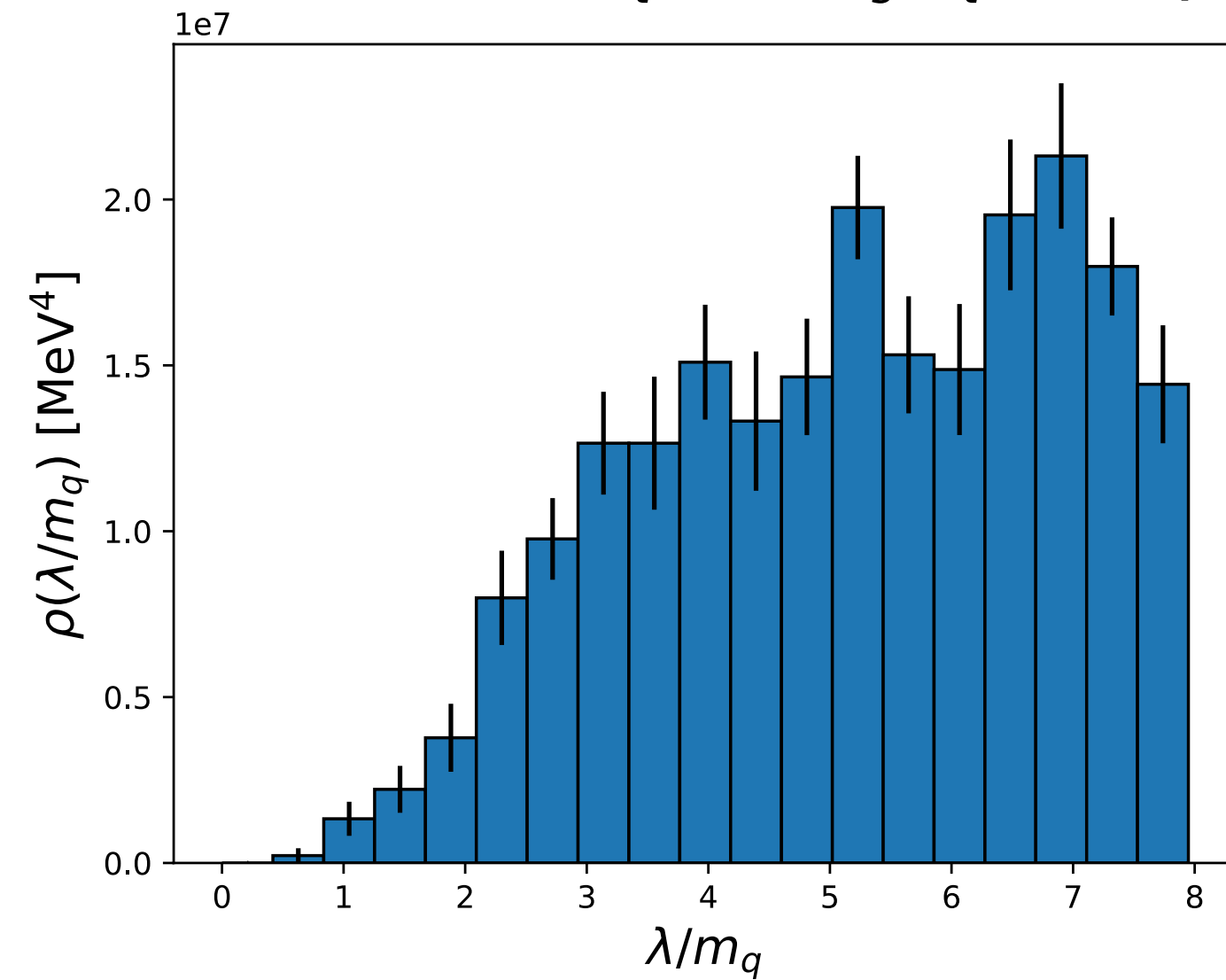
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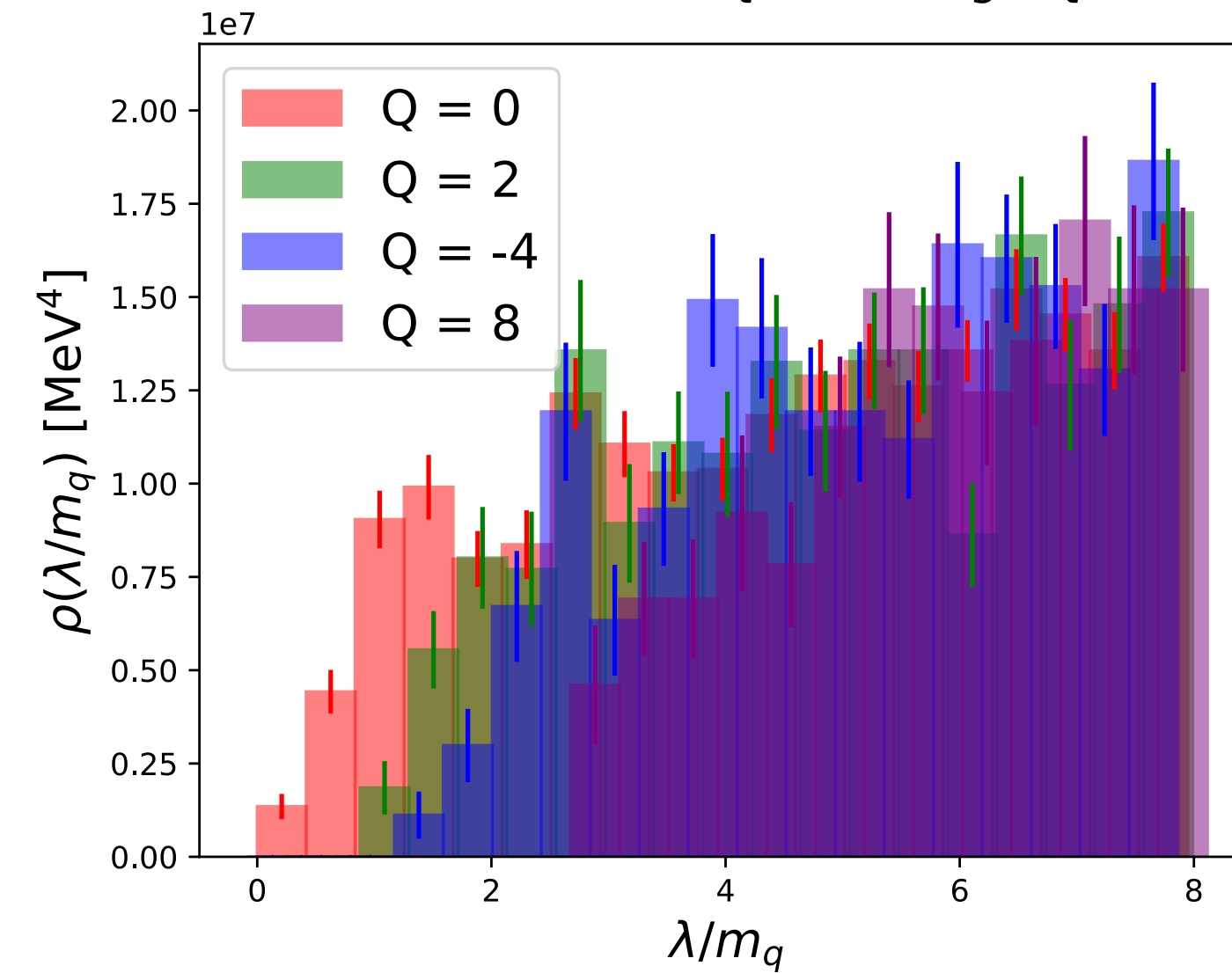
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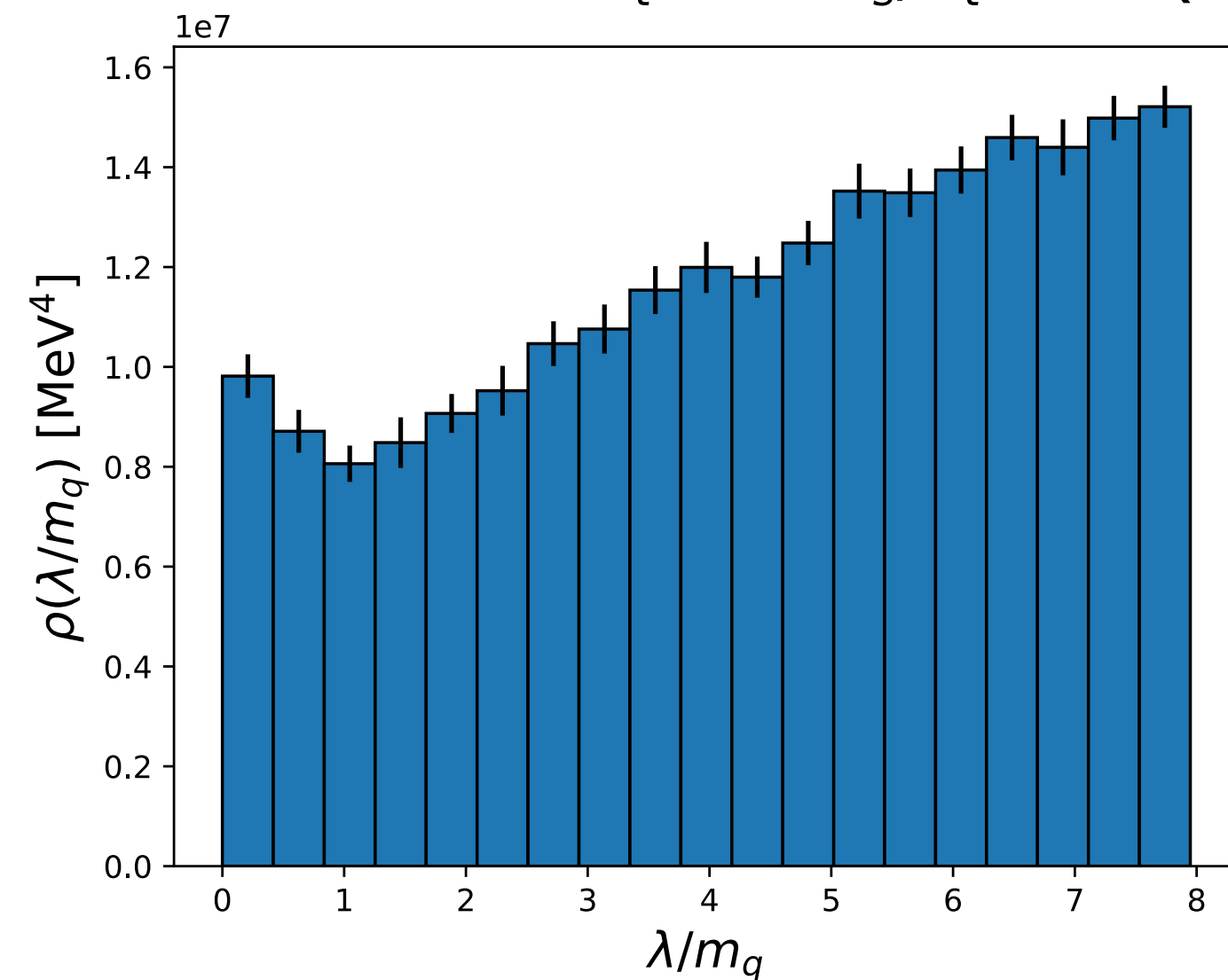
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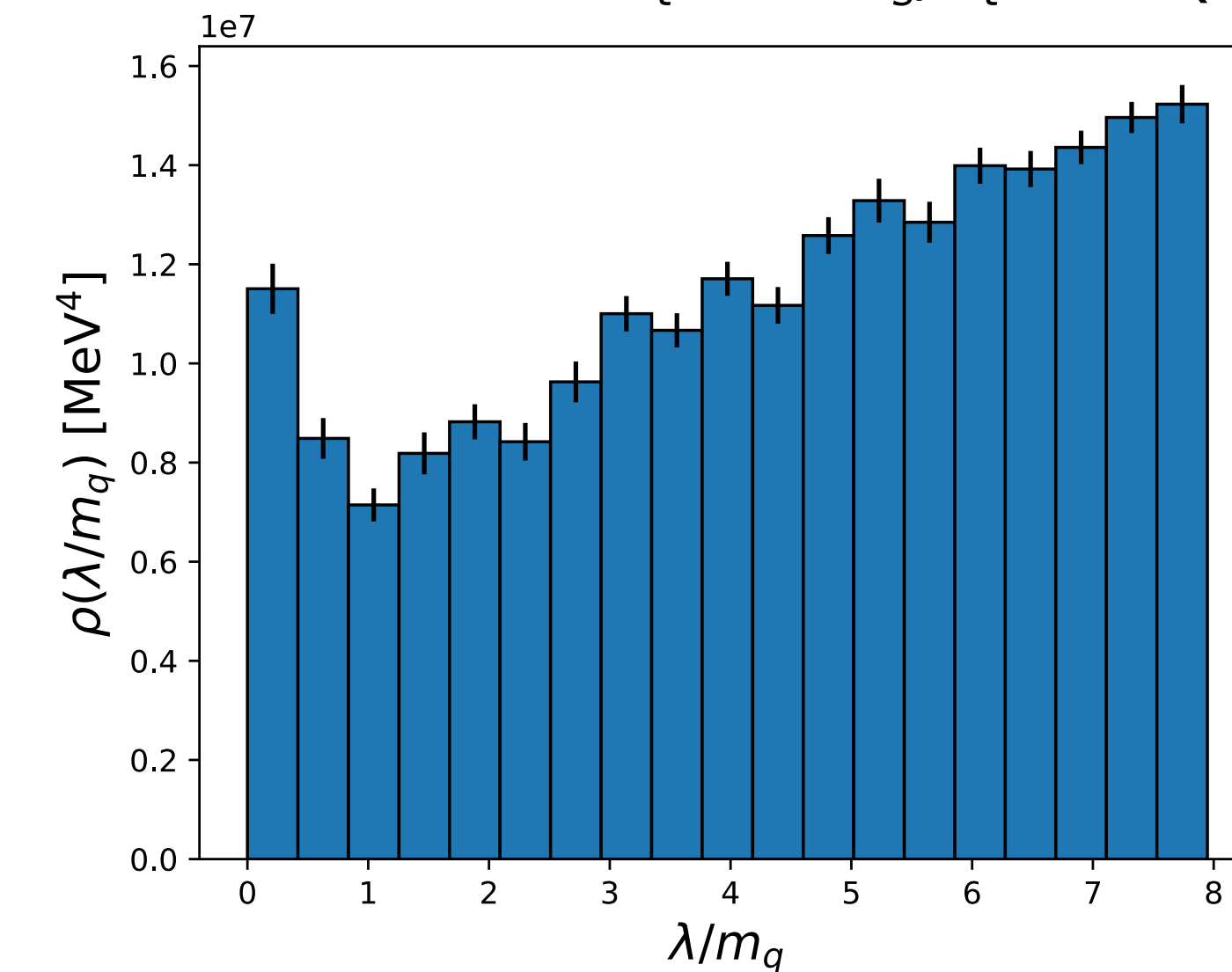
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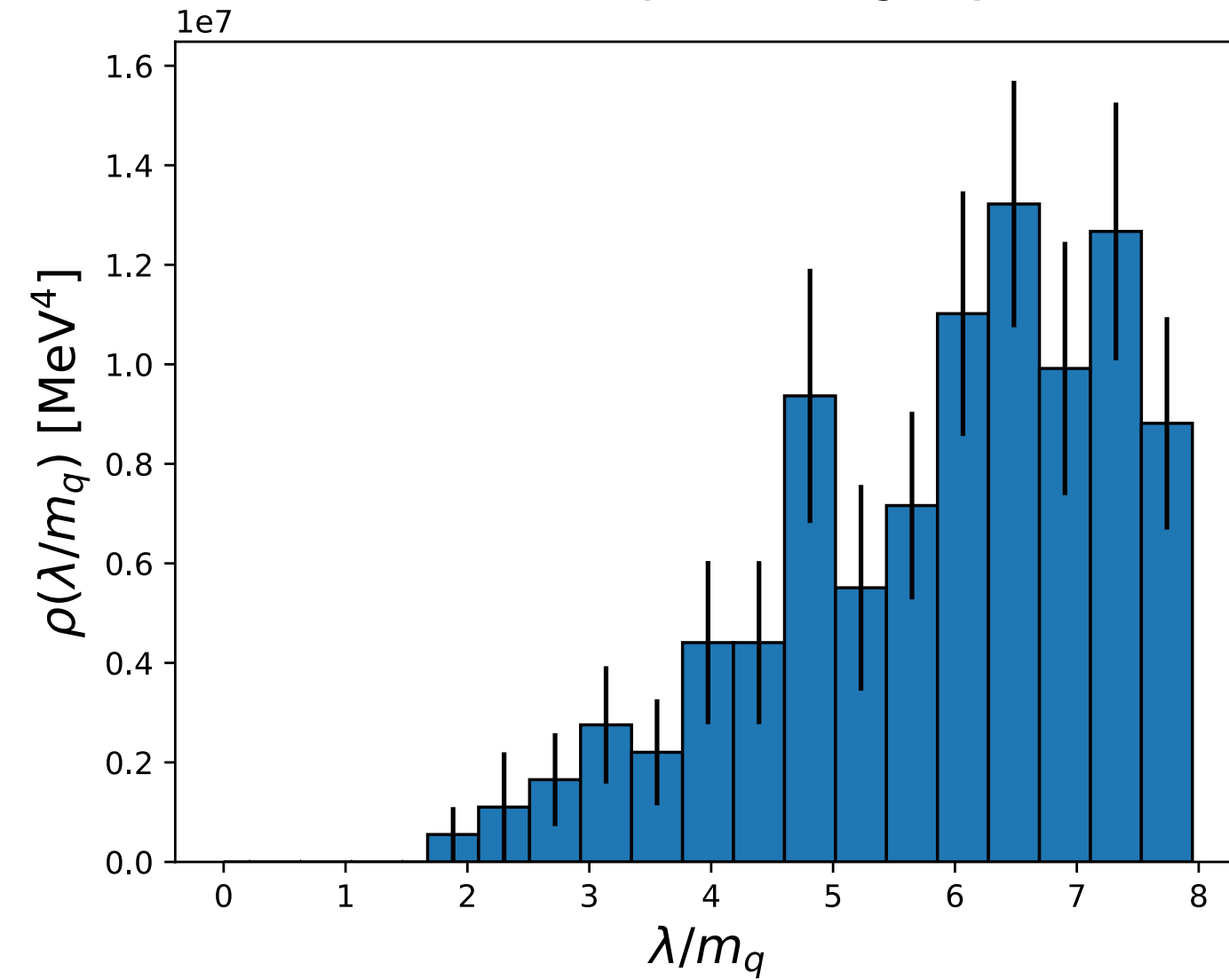
$T = 145.0 \text{ MeV}$   $N_t = 8$   $N_s/N_t = 5$   $Q = 0$



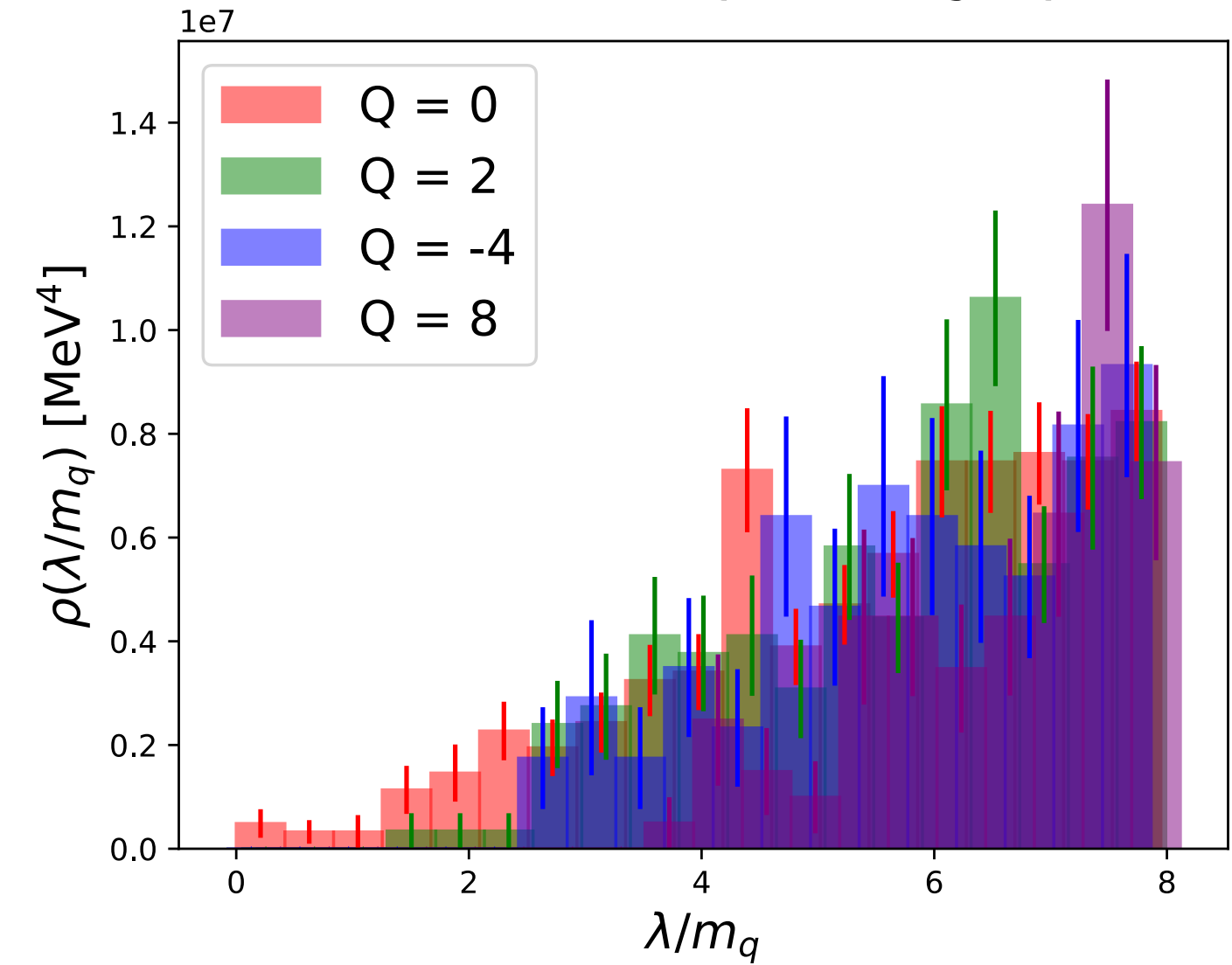
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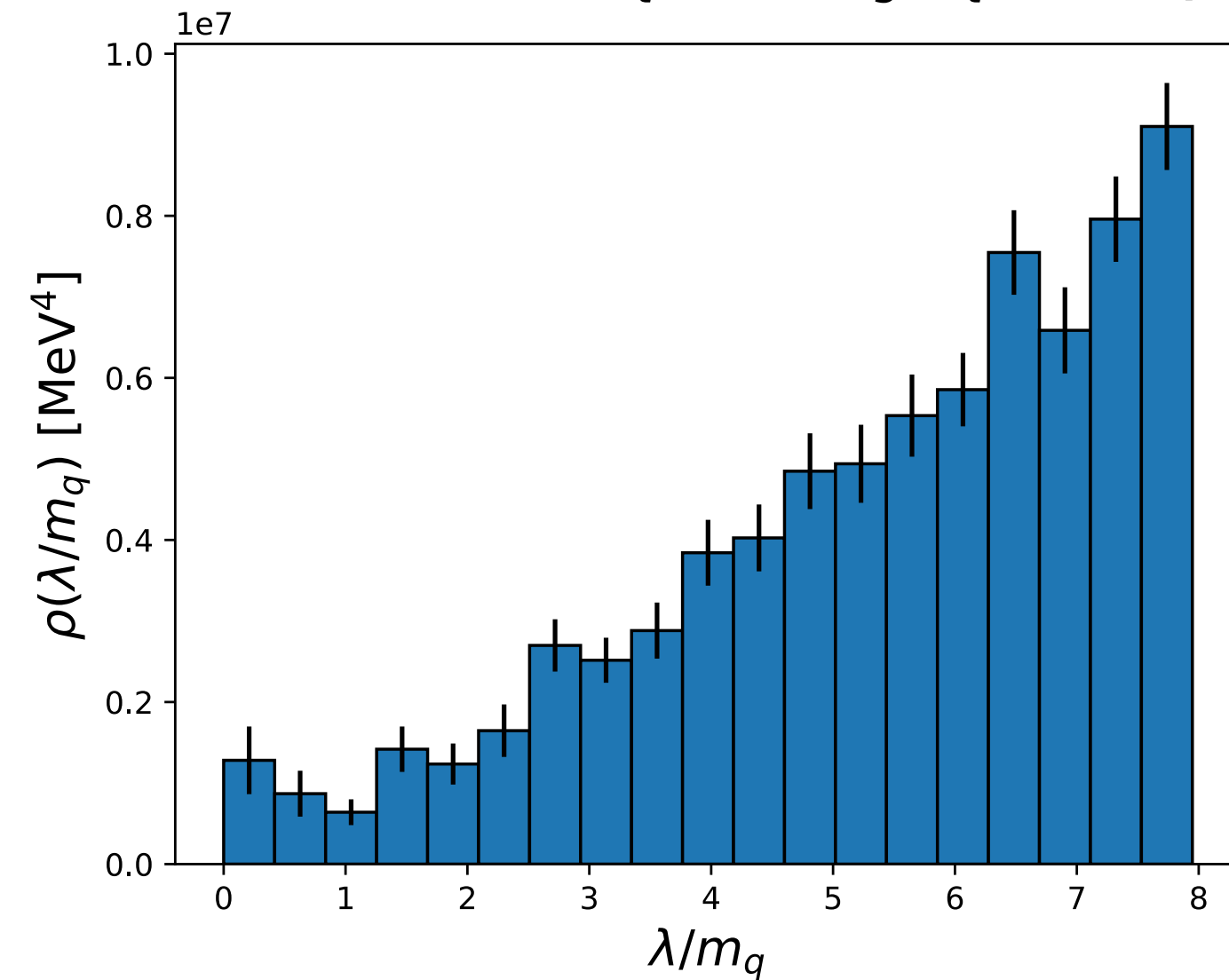
$T = 170.0 \text{ MeV}$   $N_t = 8$   $N_s/N_t = 2$   $Q = 0$



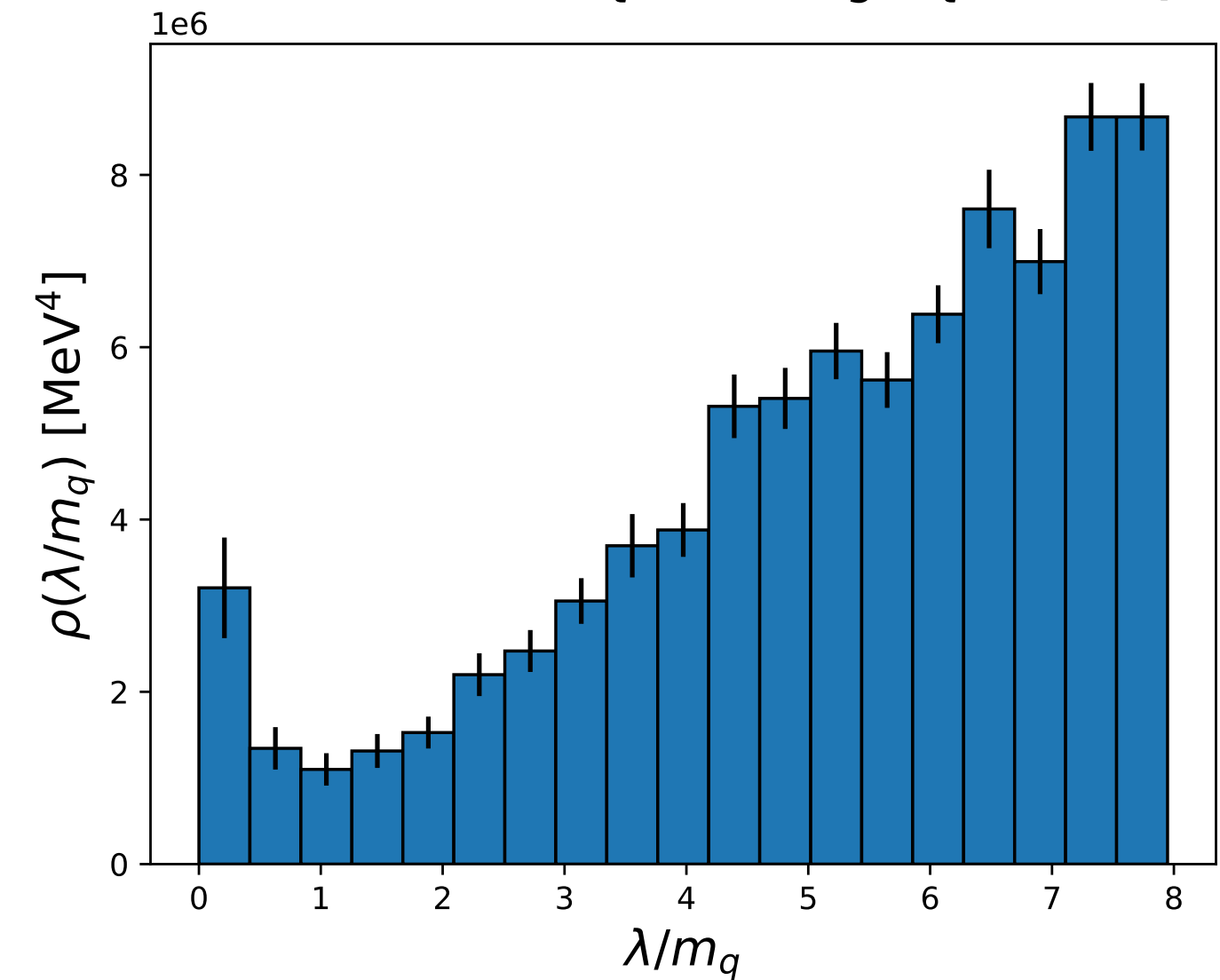
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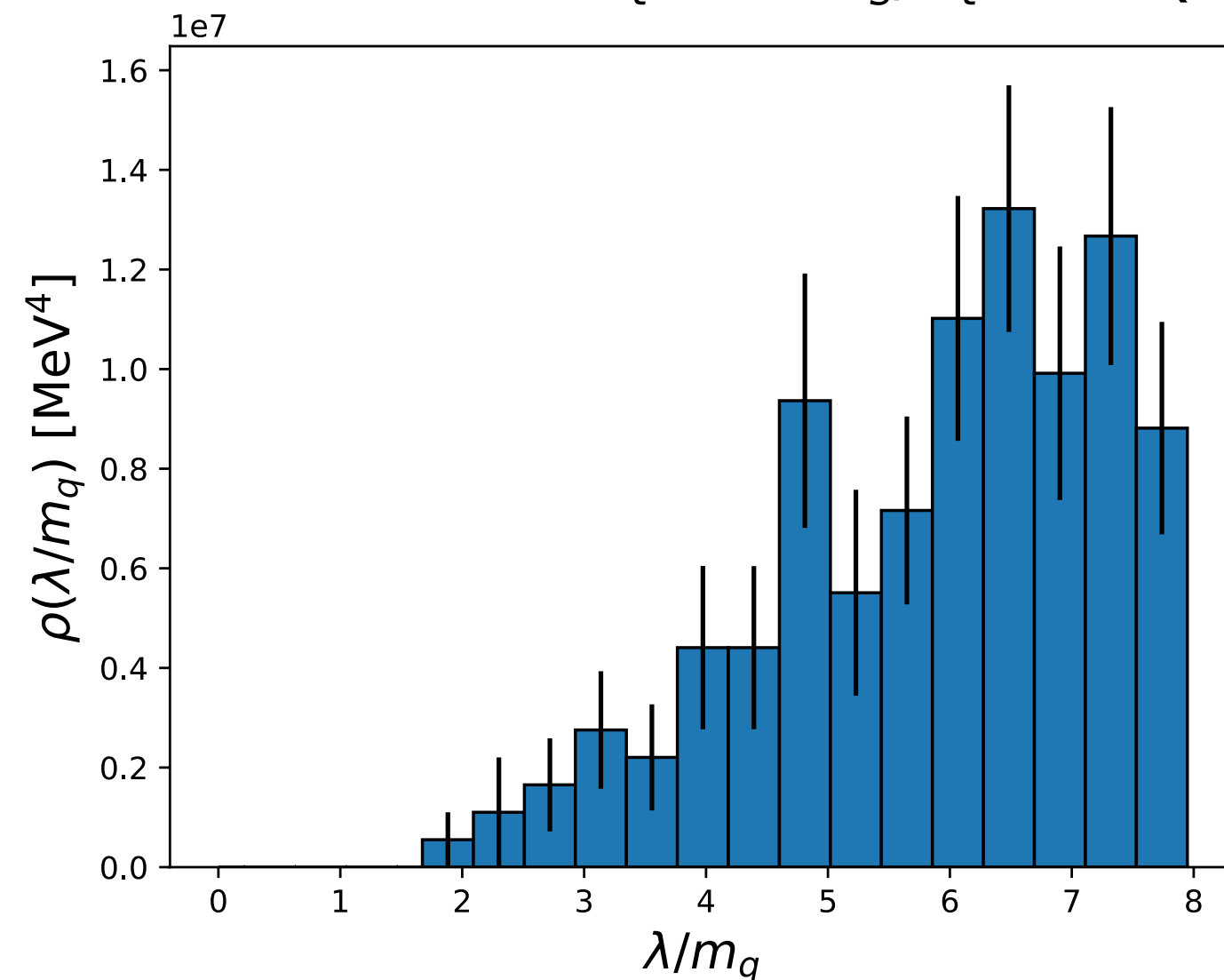
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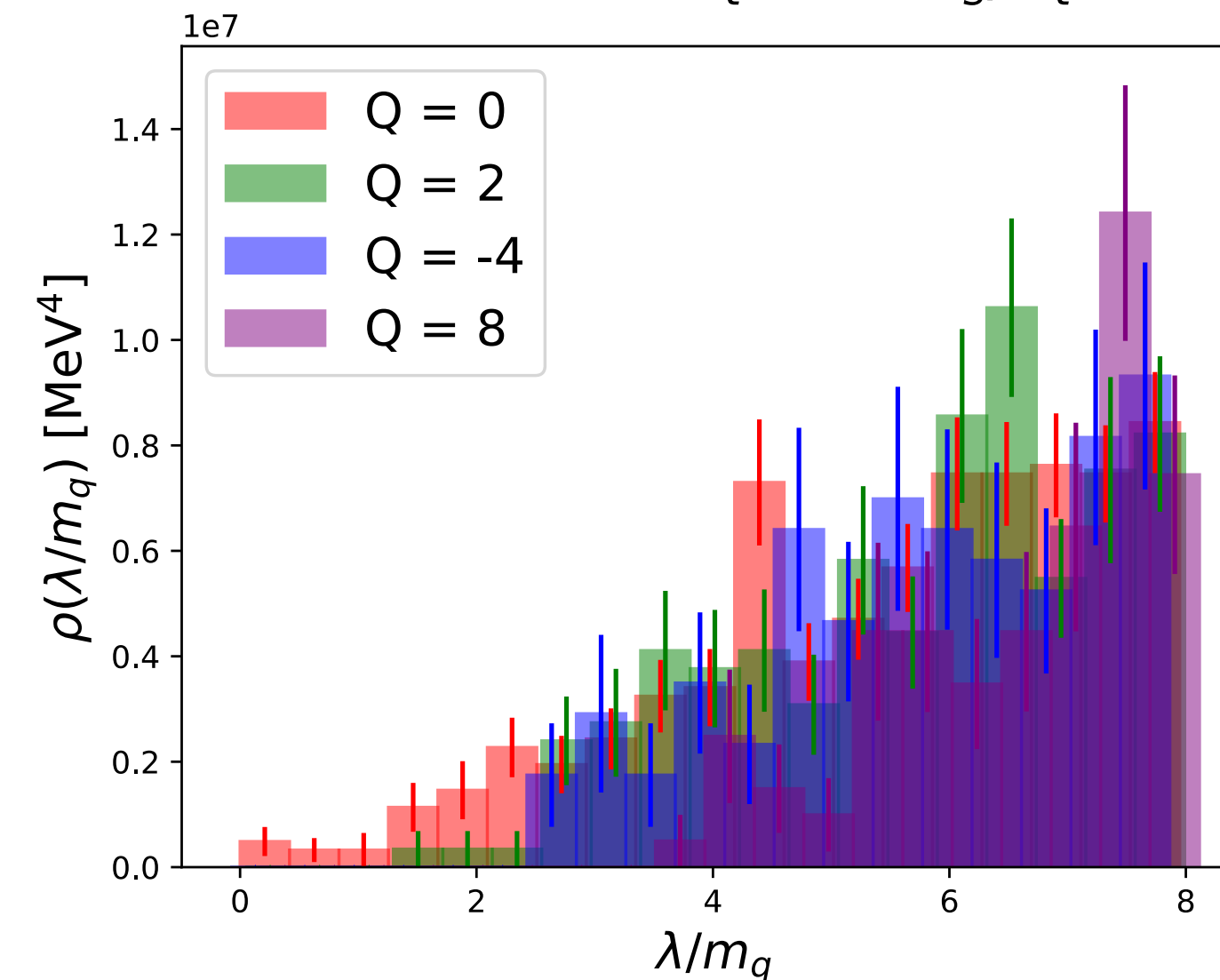
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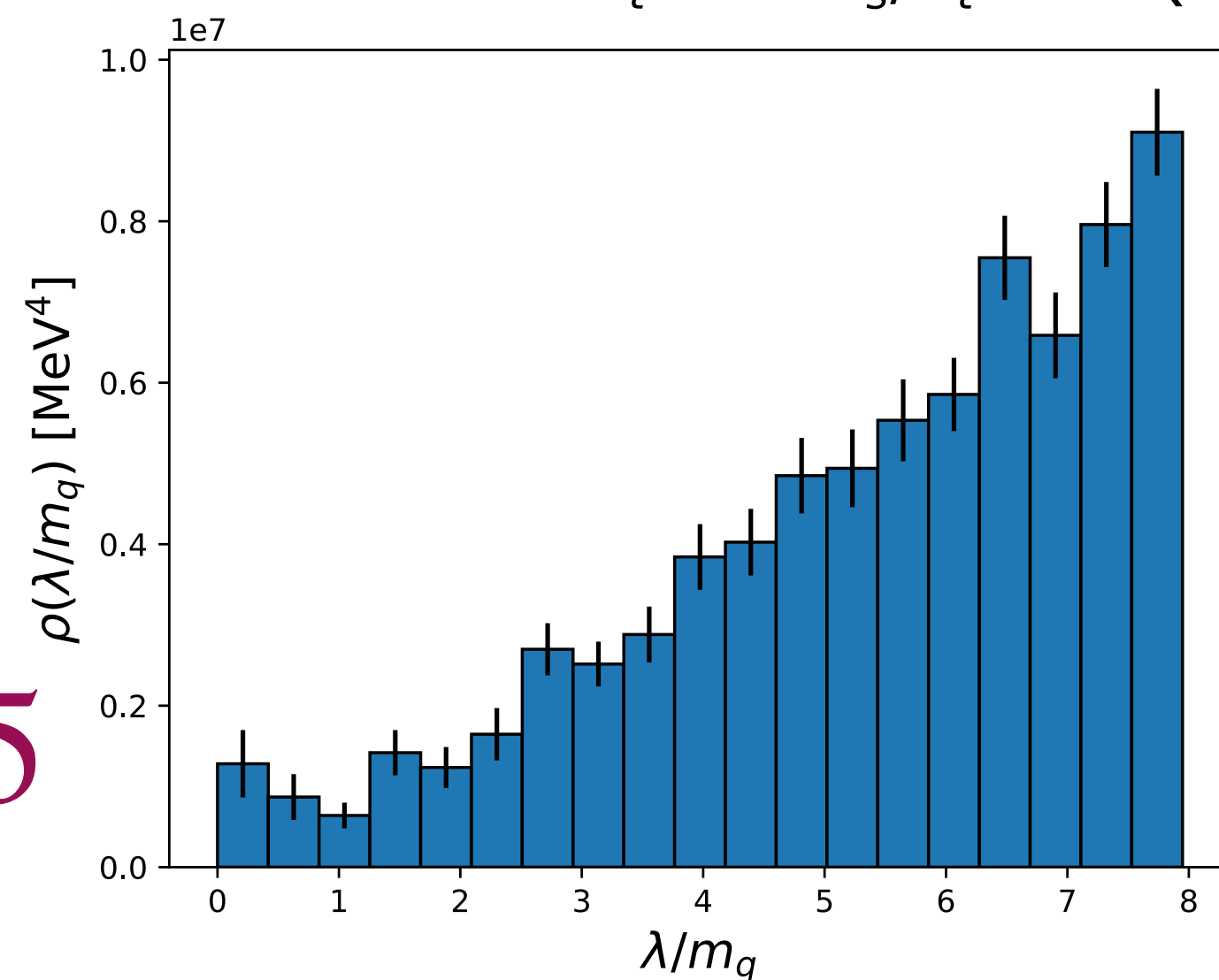
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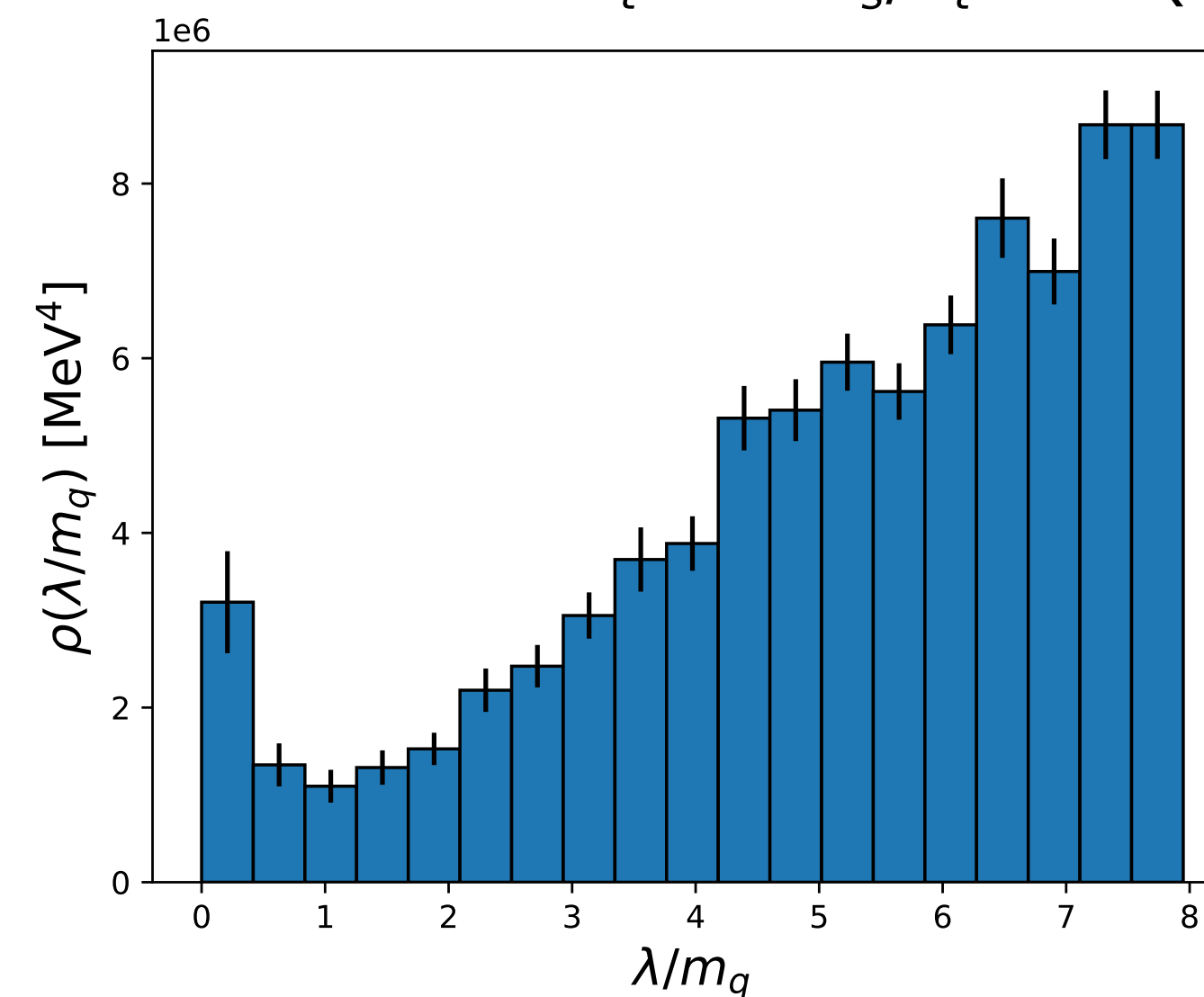
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Peak  $\rho(\lambda \rightarrow 0)$ :  
Large  $N_s/N_t \gtrsim 4 - 5$

# Localisation

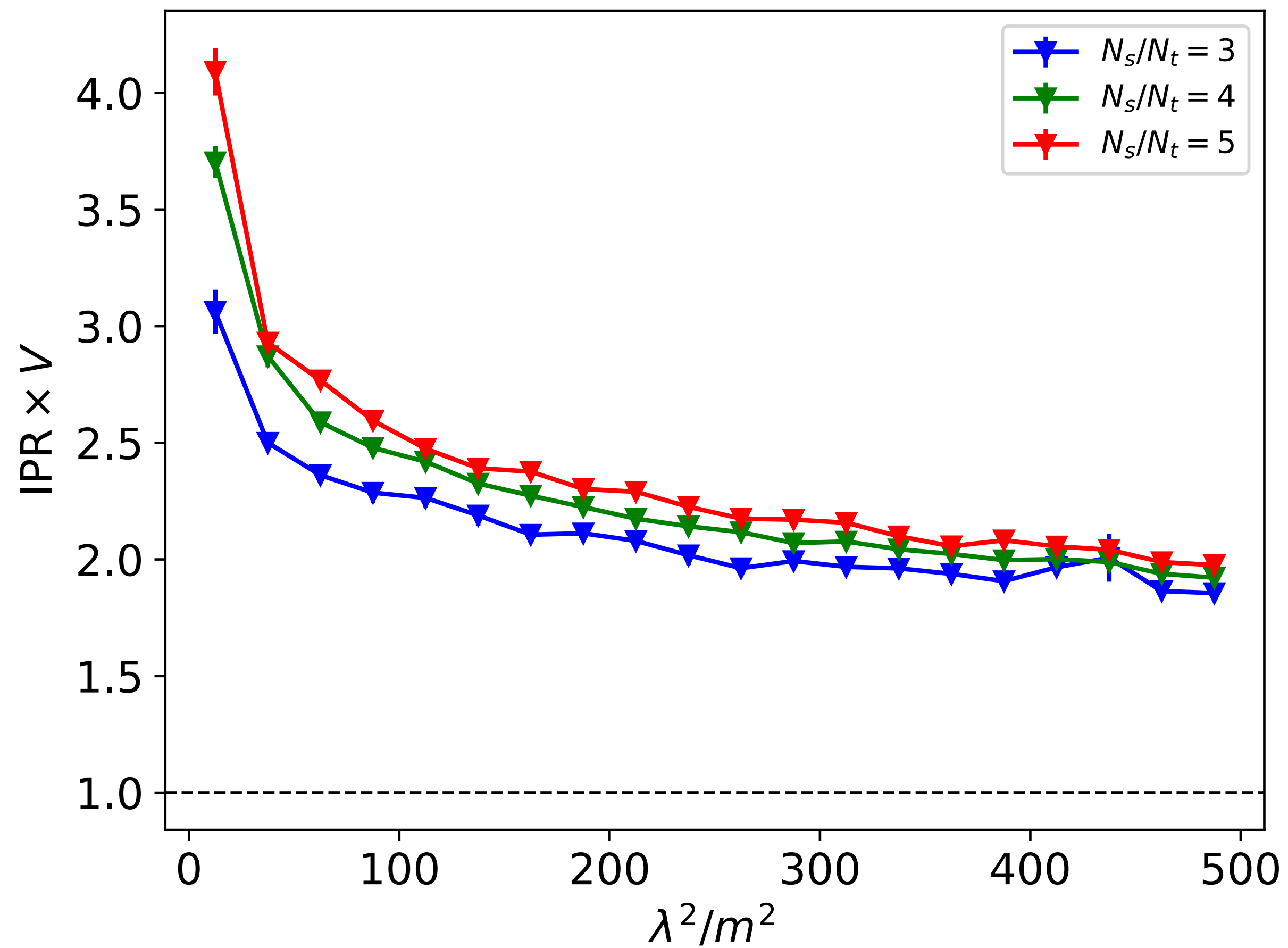
## Eigenmodes of $D$

- $\text{IPR}_n = \sum_x ||\psi_n(x)||^4$
- $\text{IPR}(\lambda) = \frac{\langle \sum_n \delta(\lambda - \lambda_n) \text{IPR}_n \rangle}{\langle \sum_n \delta(\lambda - \lambda_n) \rangle}$ 
  - Localised:  $\text{IPR} \times V \sim V$
  - Delocalised:  $\text{IPR} \times V \sim \text{const}$

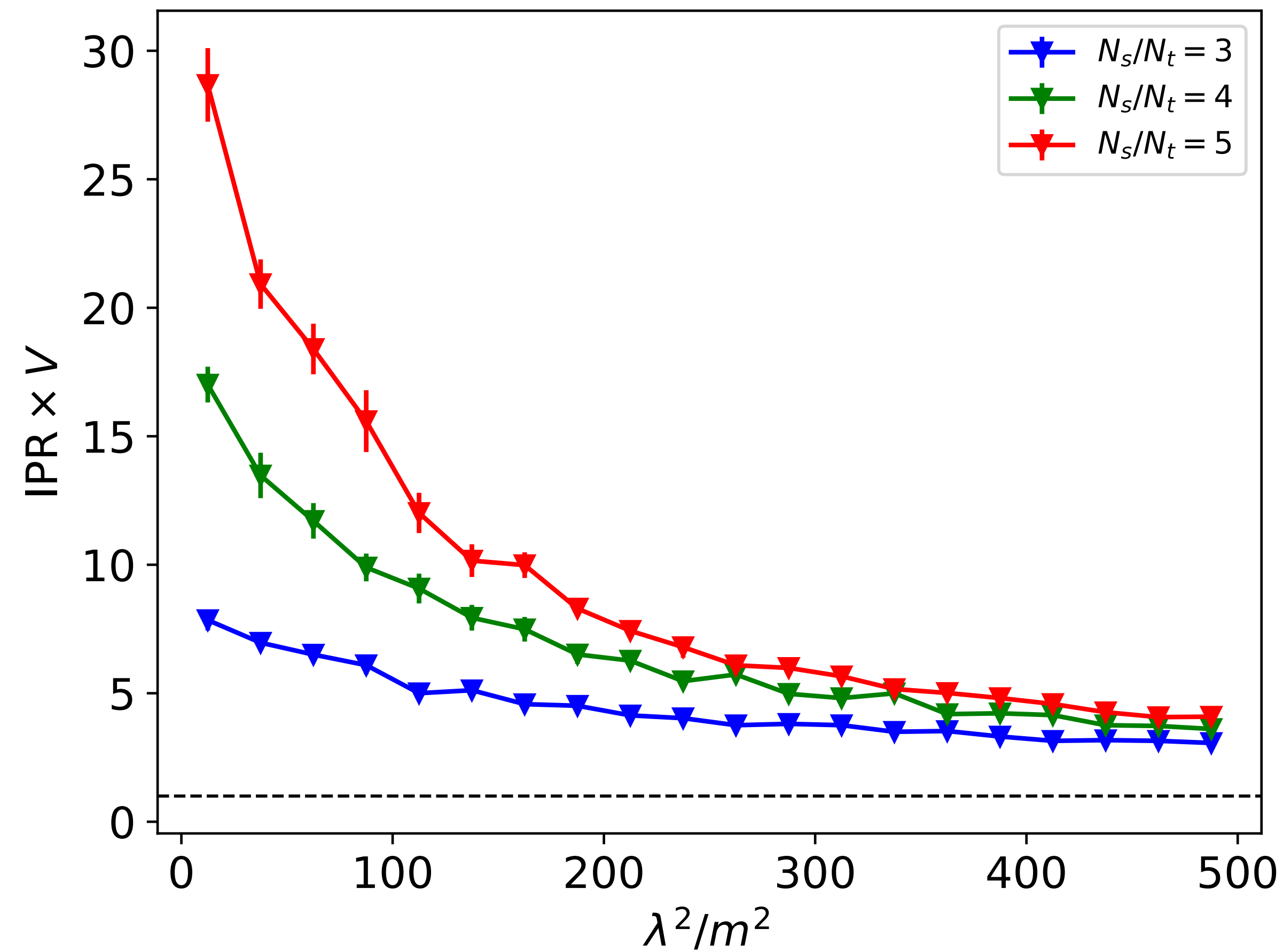
# Localisation

$$N_t = 8$$

T=135 [MeV]



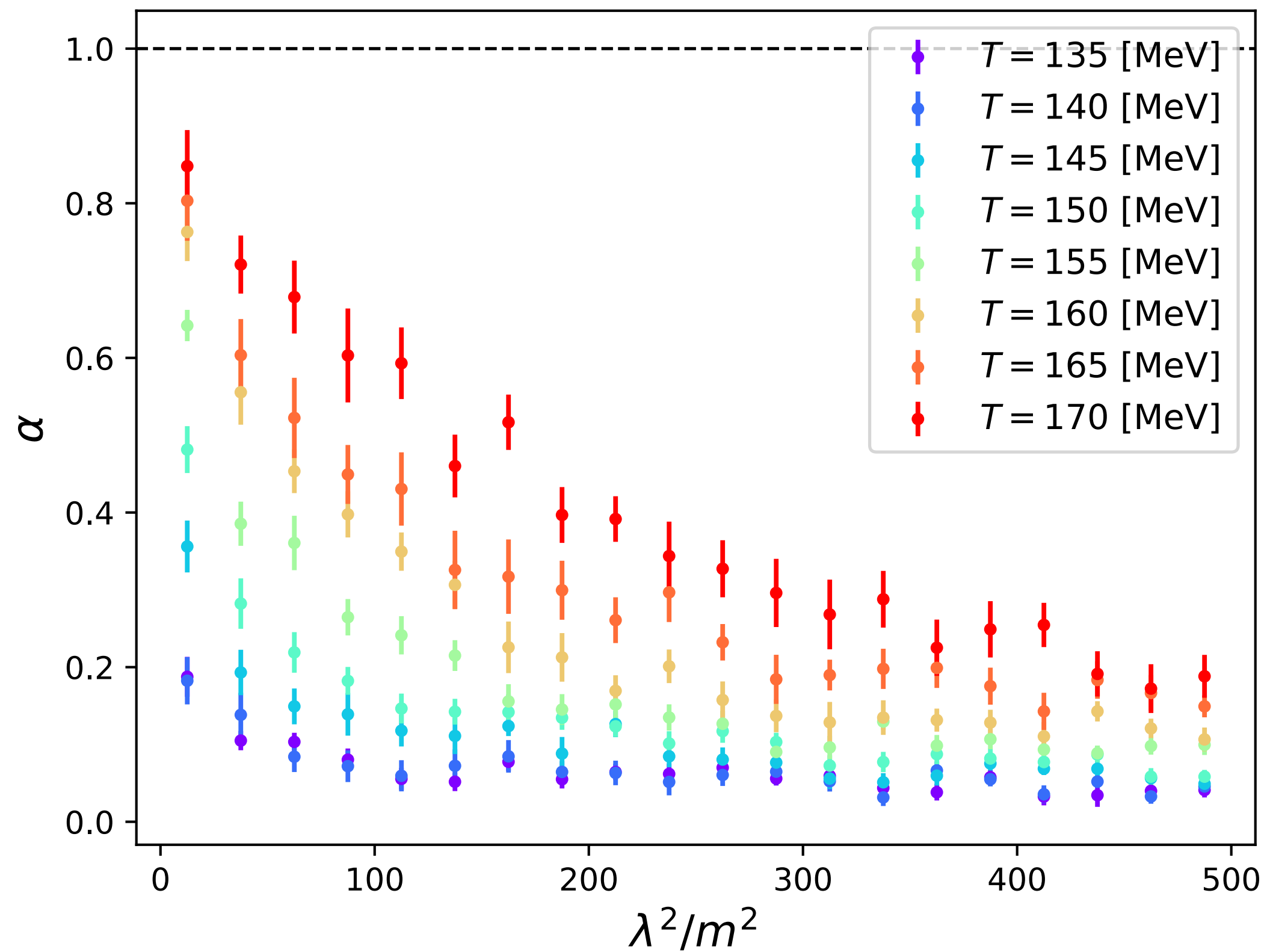
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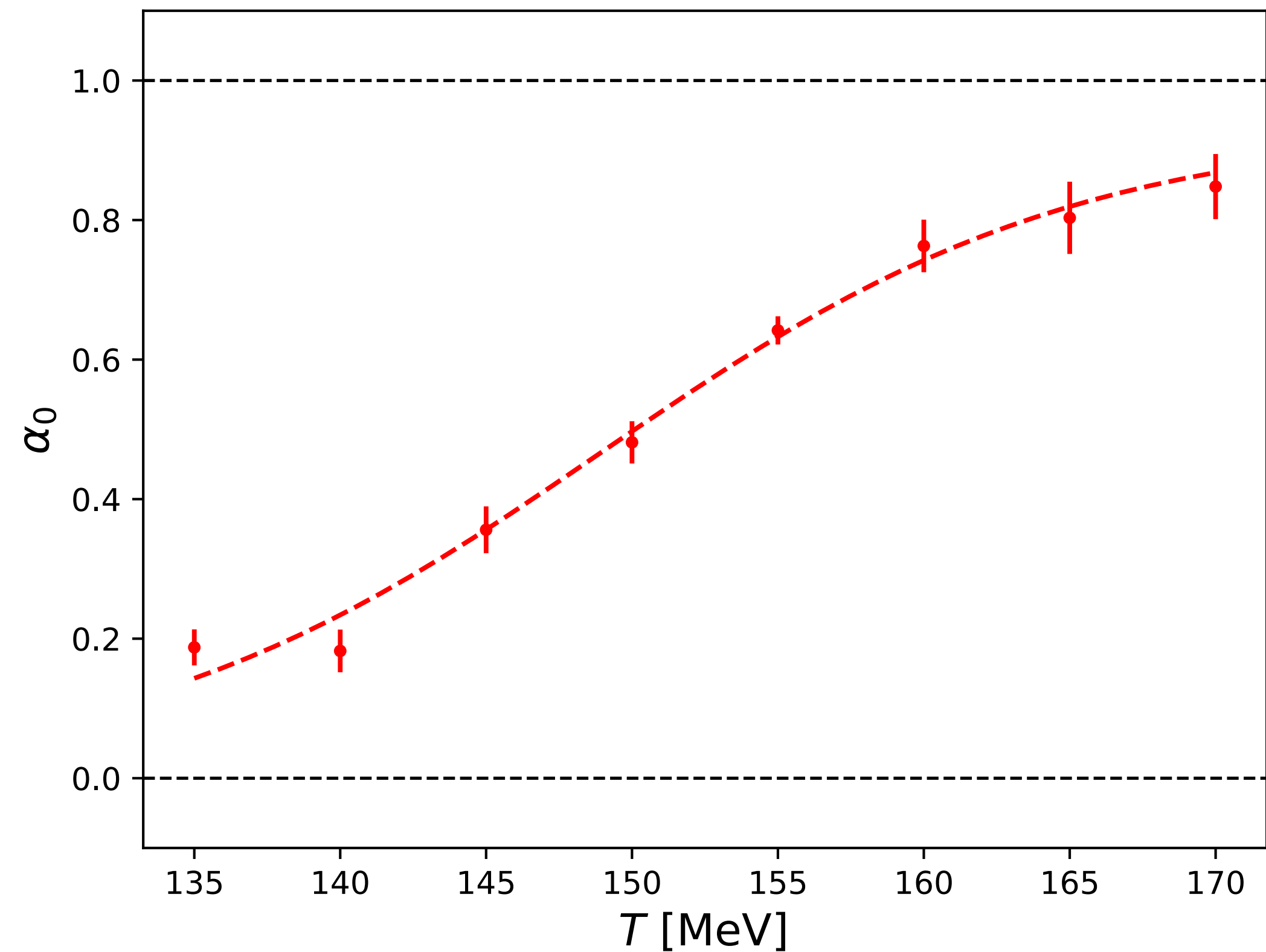
# Volume scaling

$$\text{IPR} \times V \sim V^\alpha$$

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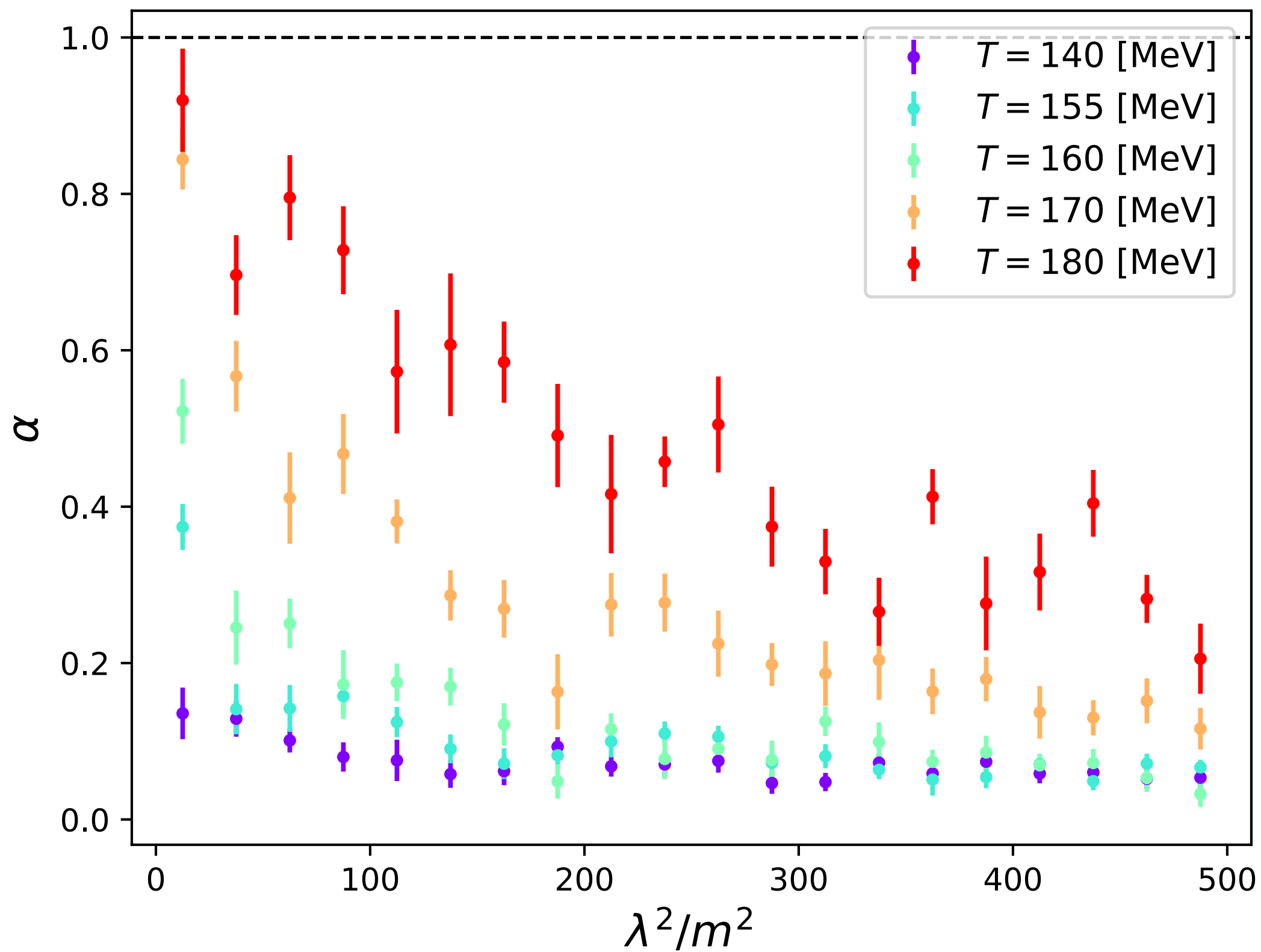


$T_l(N_t = 8) \approx 150$  MeV: consistent with  $T_c(N_t = 8)$

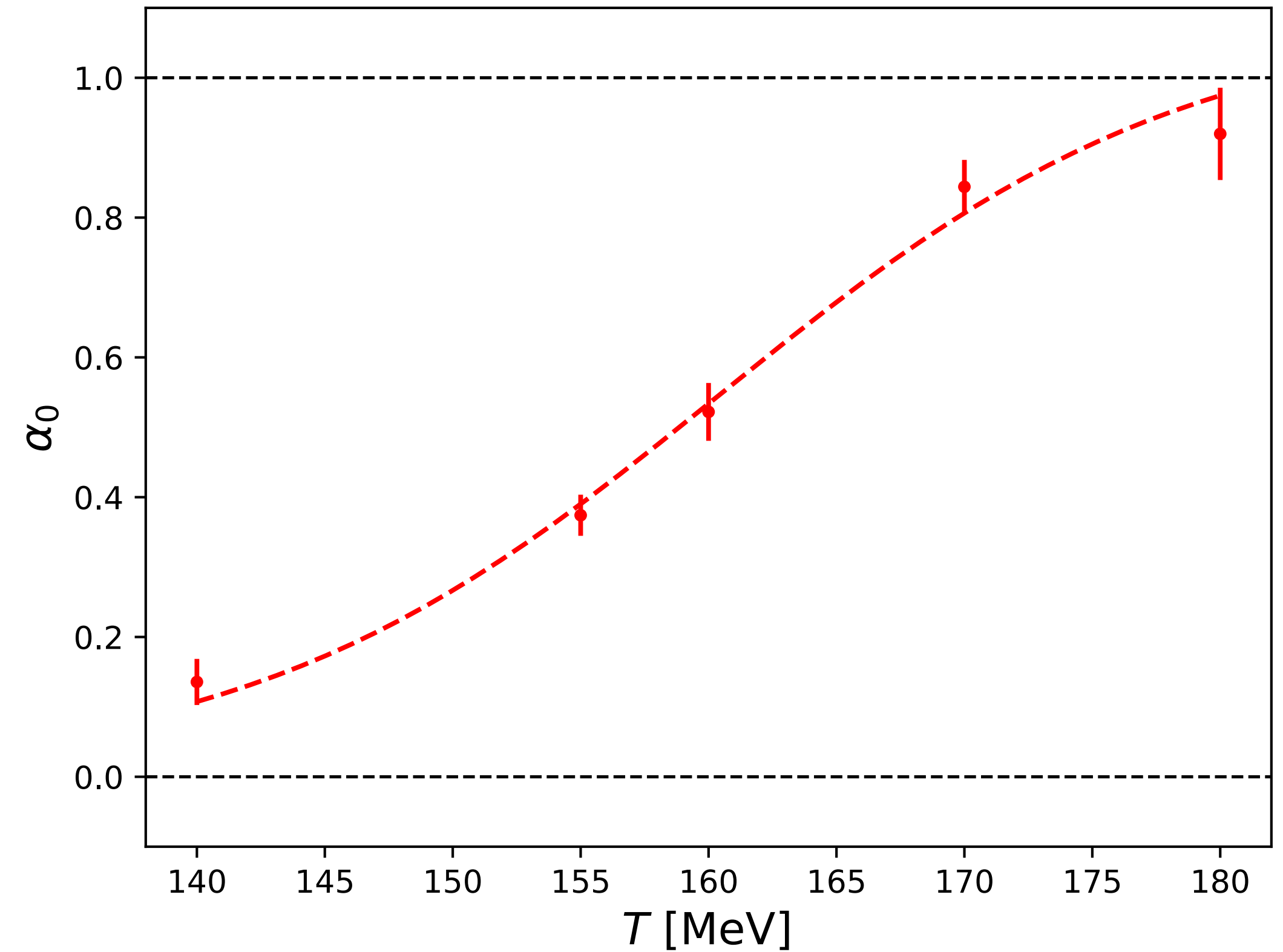
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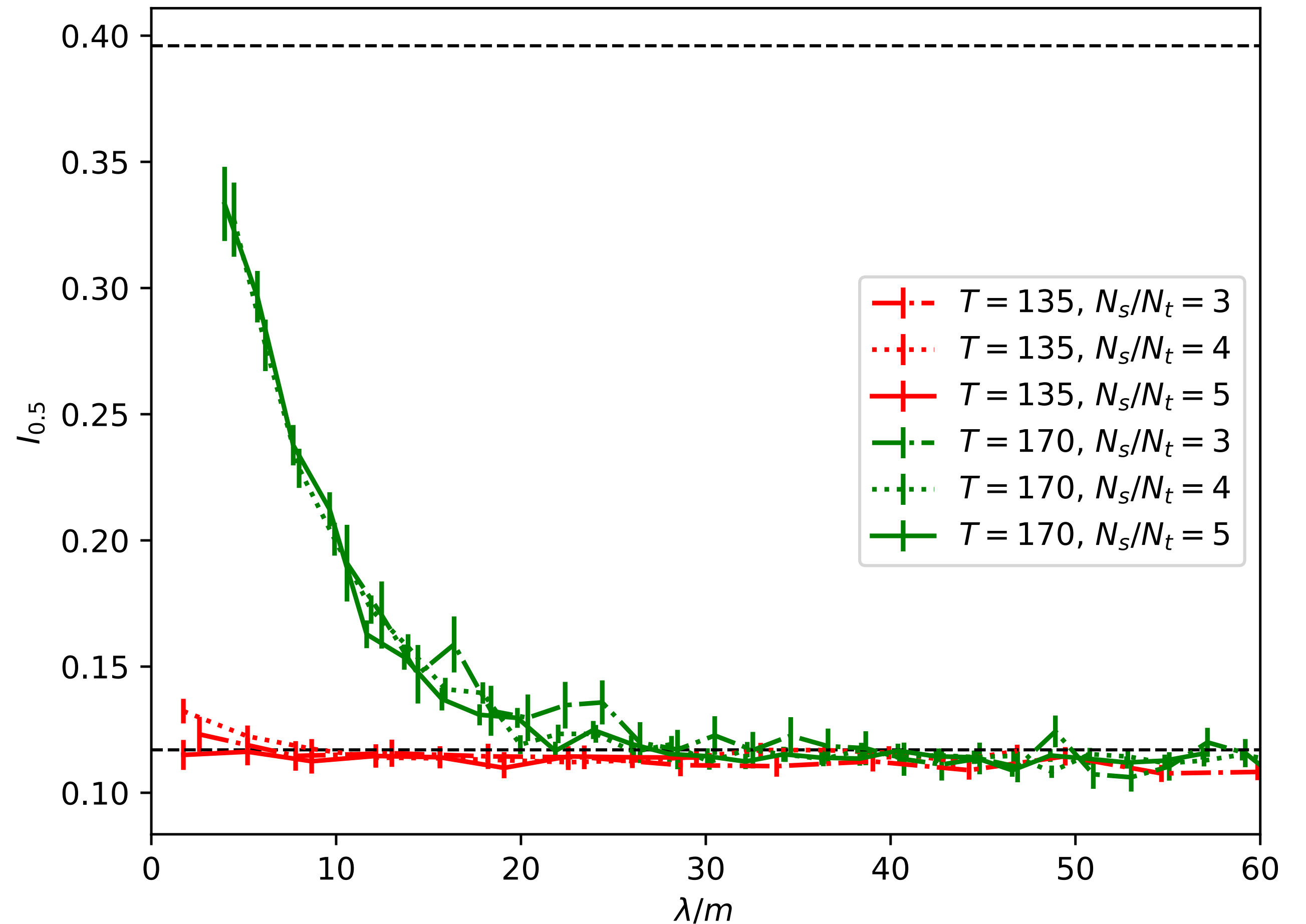
$T_l(N_t = 10) \approx 160$  MeV: consistent with  $T_c(N_t = 10)$

# Unfolded level spacing distribution

- Unfolding  $x_i = \int^{\lambda_i} d\lambda \rho(\lambda)$      $s_i = x_{i+1} - x_i$
- Unfolded level spacing distribution (ULSD):

$$p(s, \lambda) = \frac{\langle \sum_n \delta(\lambda - \lambda_n) \delta(s - s_n) \rangle}{\rho(\lambda)}$$

- Integrated ULSD  $I_{s_0} = \int_0^{s_0} ds p(s, \lambda, N_s)$ 
  - Localised:  $I_{0.5}^{\text{Poisson}} \approx 0.396$
  - Delocalised:  $I_{0.5}^{\text{RMT}} \approx 0.117$



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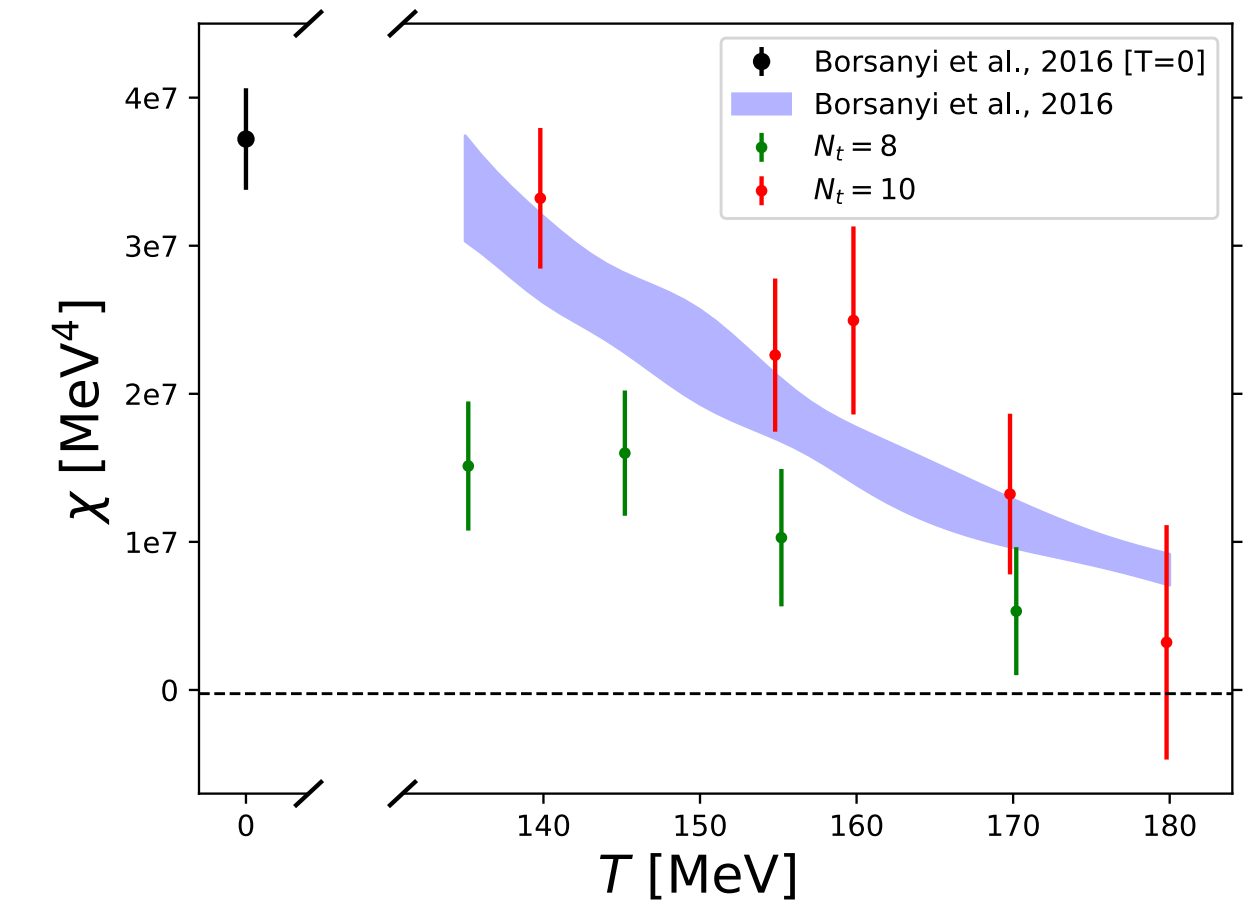
Purely overlap result!

- Dynamical overlap fermions at  $m_\pi = m_\pi^{\text{phys}}$ , summed over  $Q$

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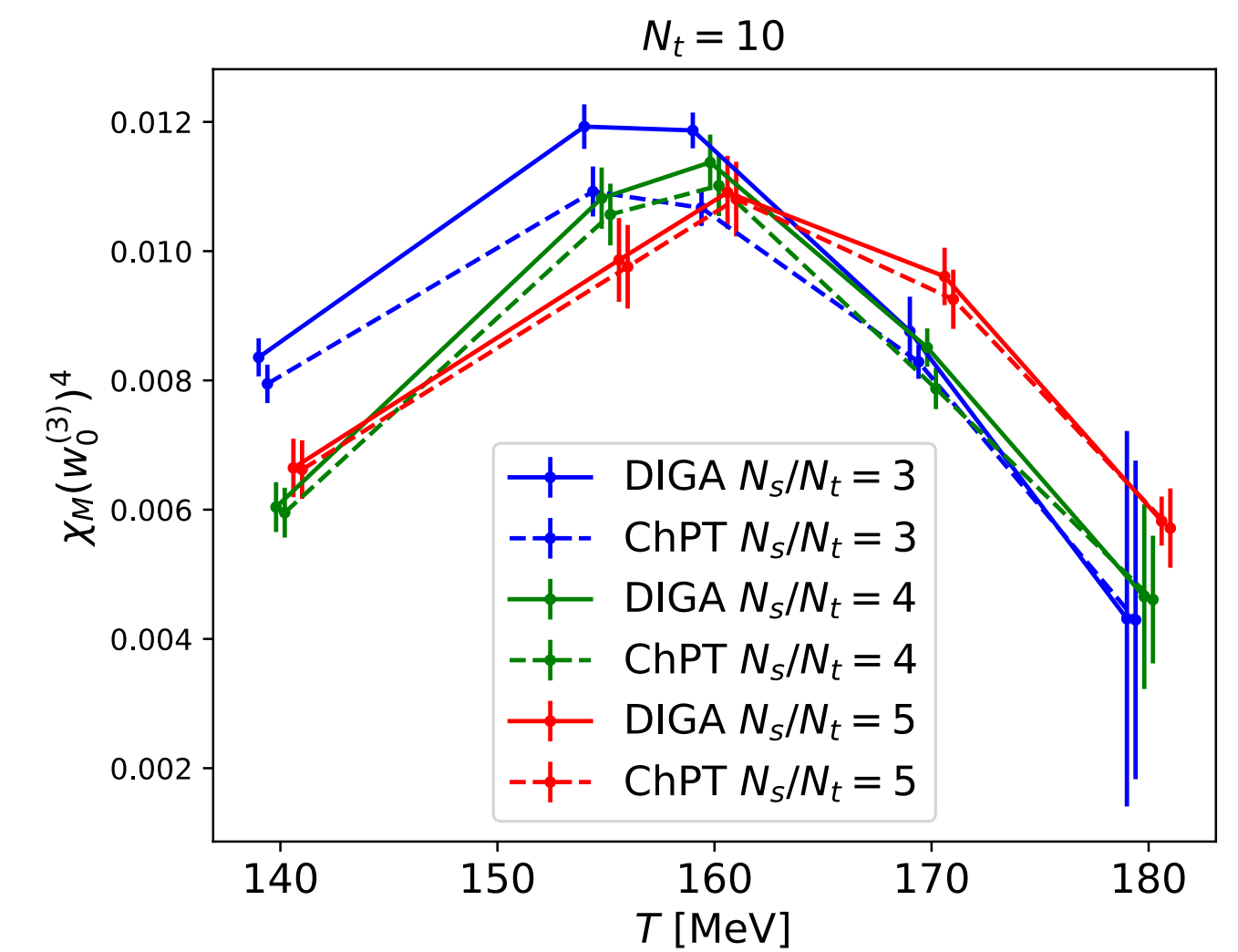
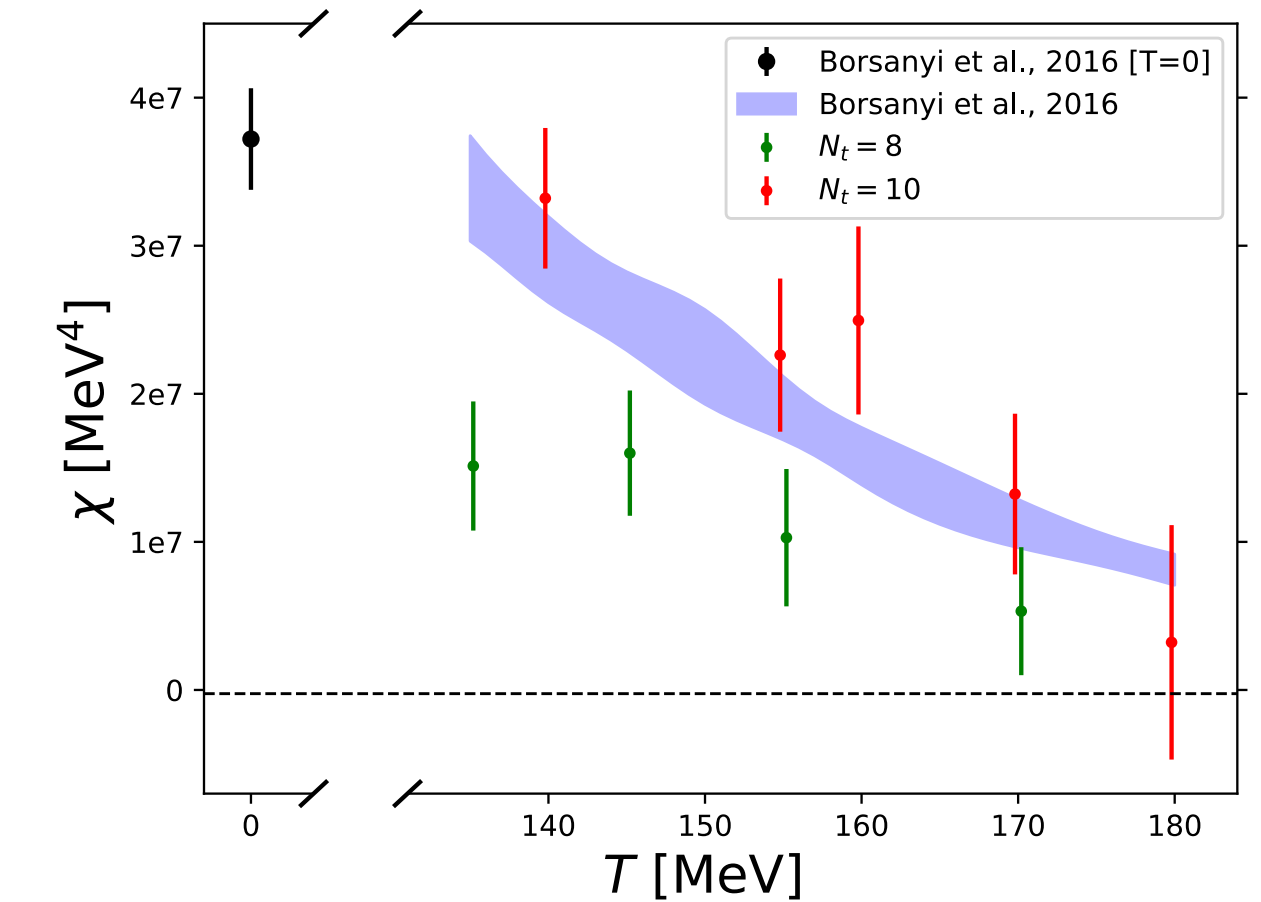
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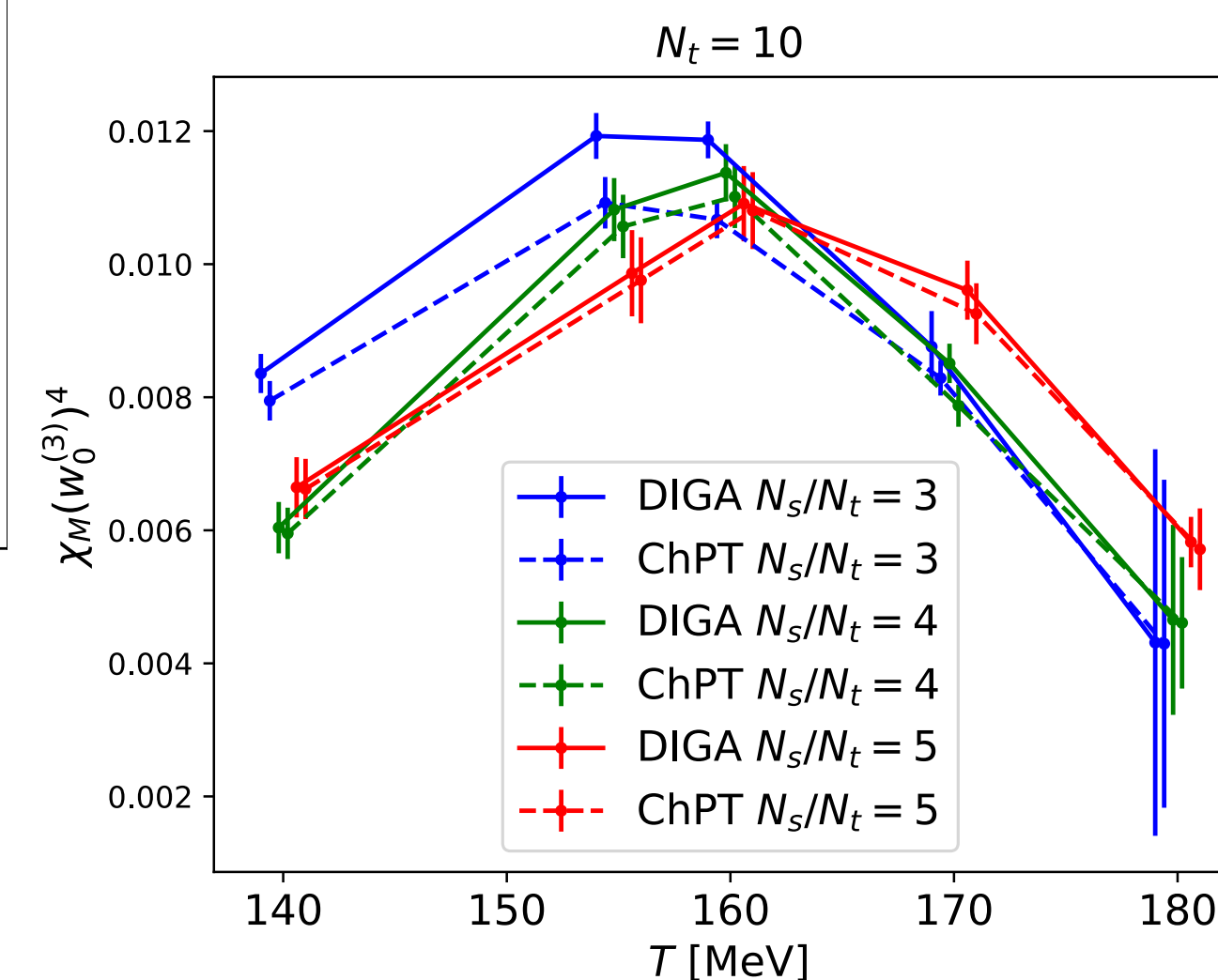
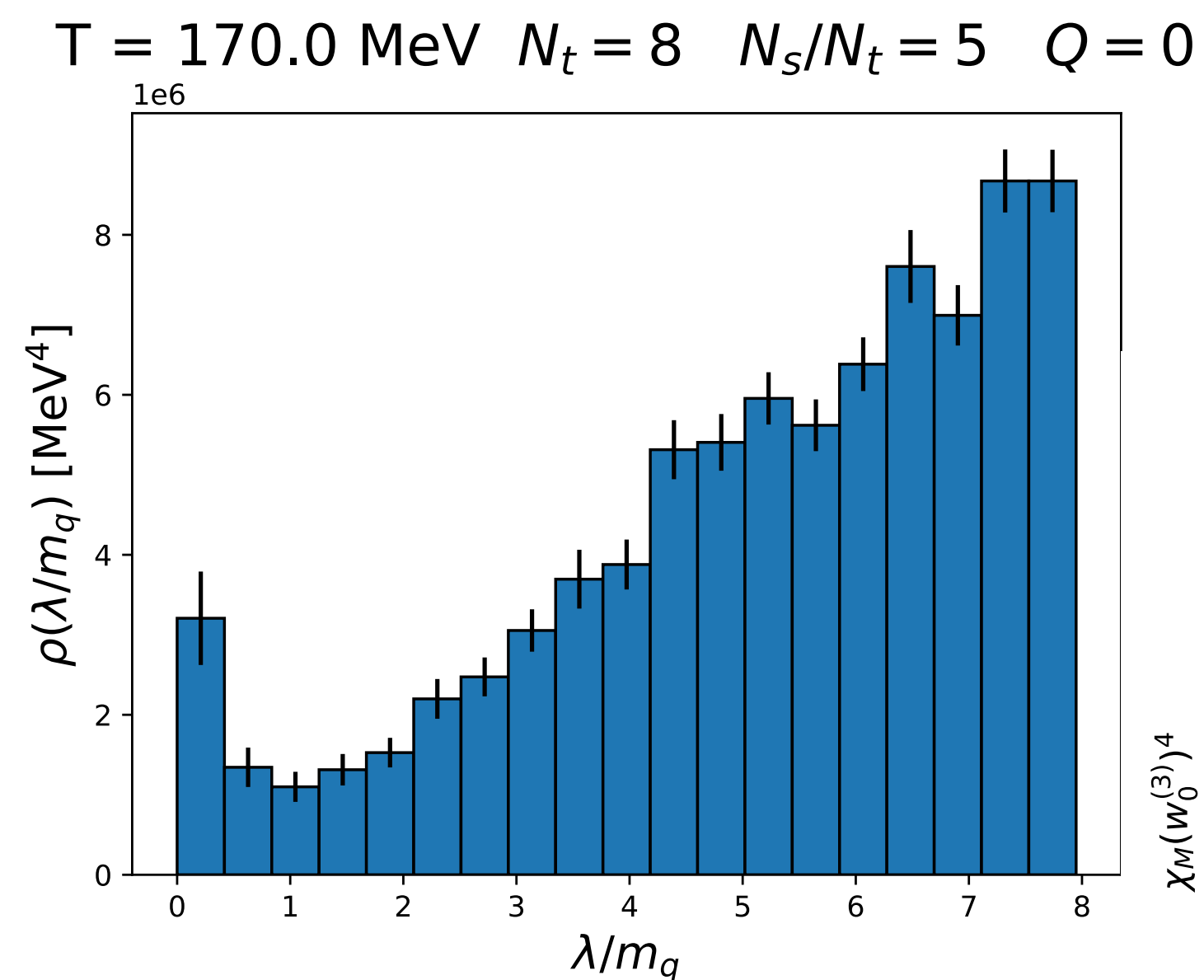
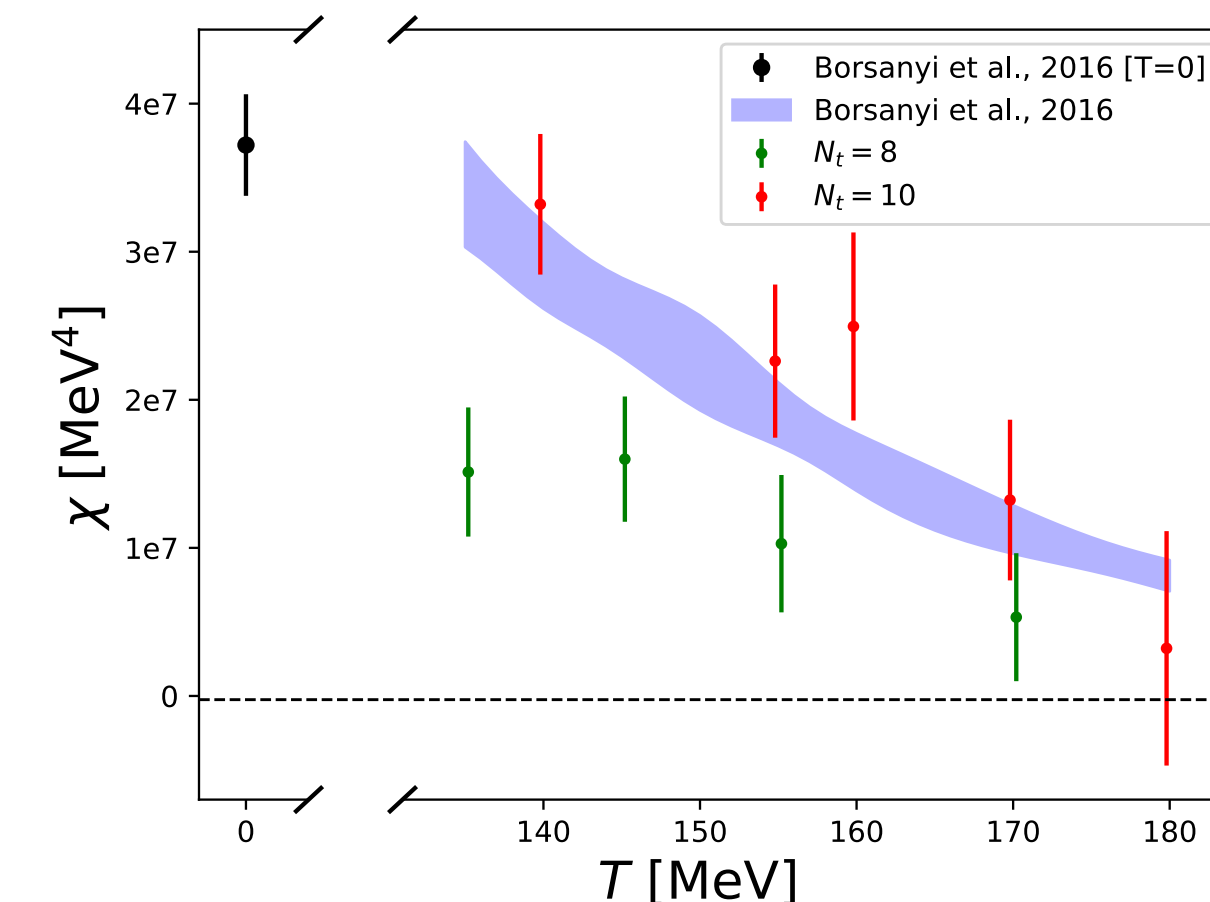
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- Consistent with chiral crossover



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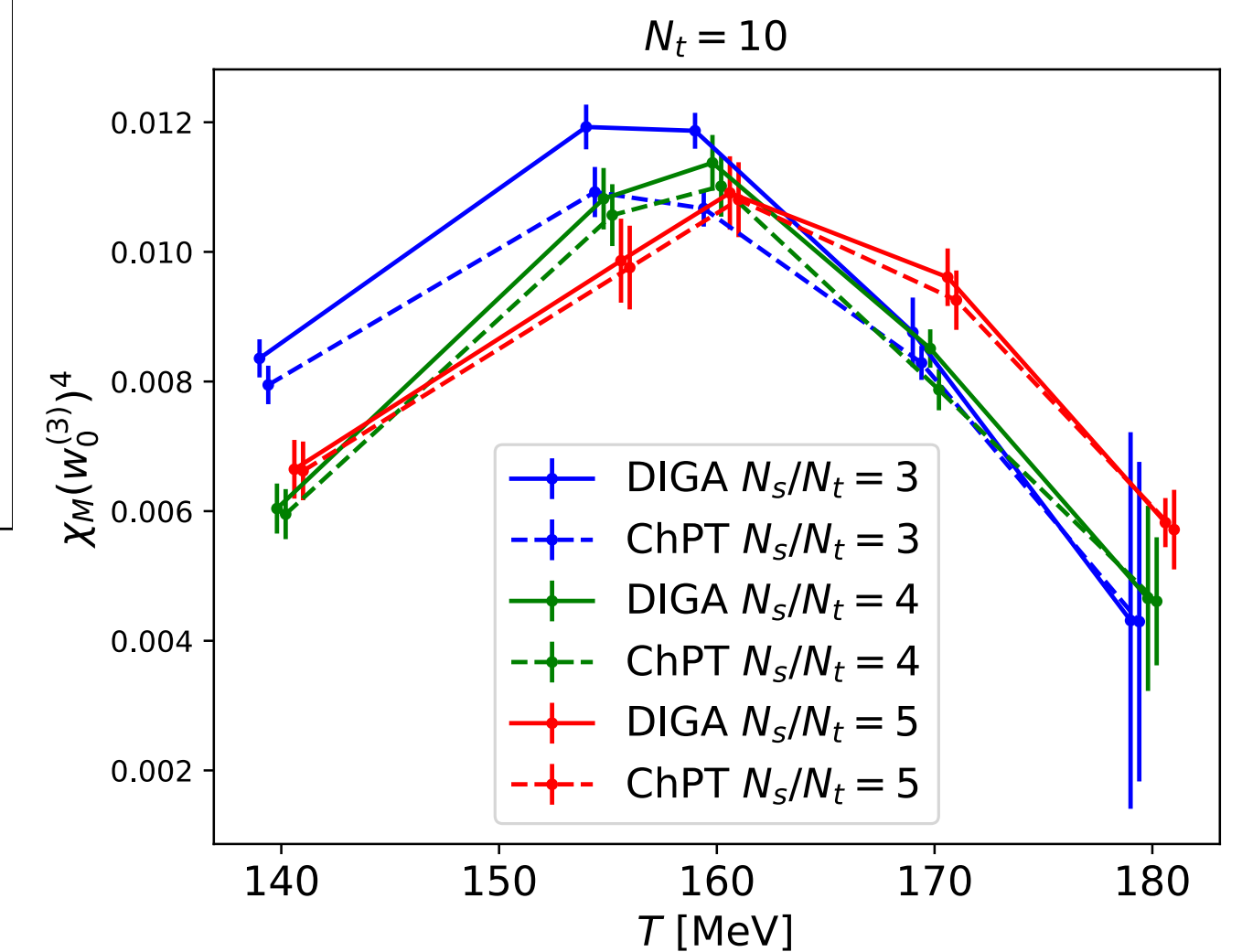
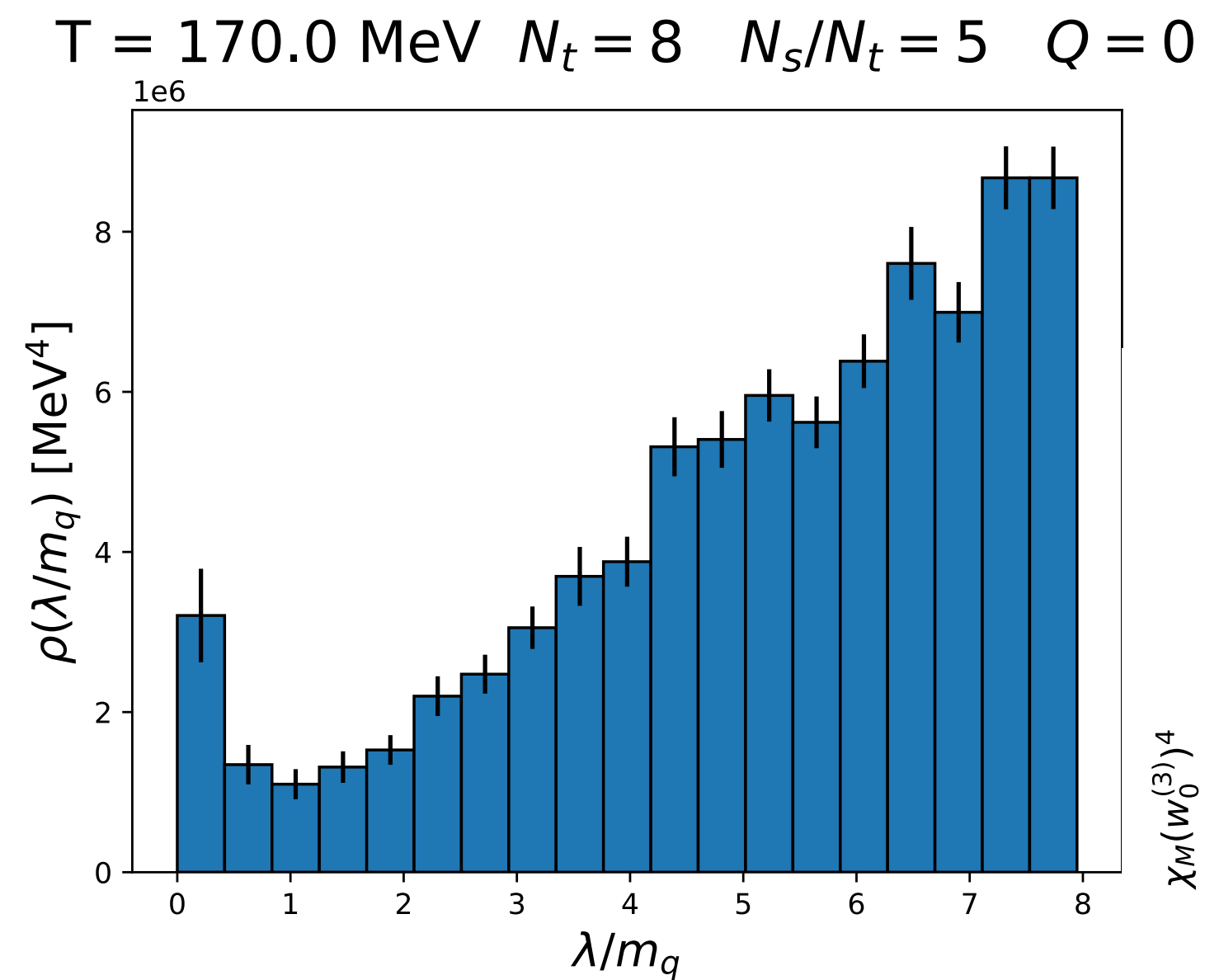
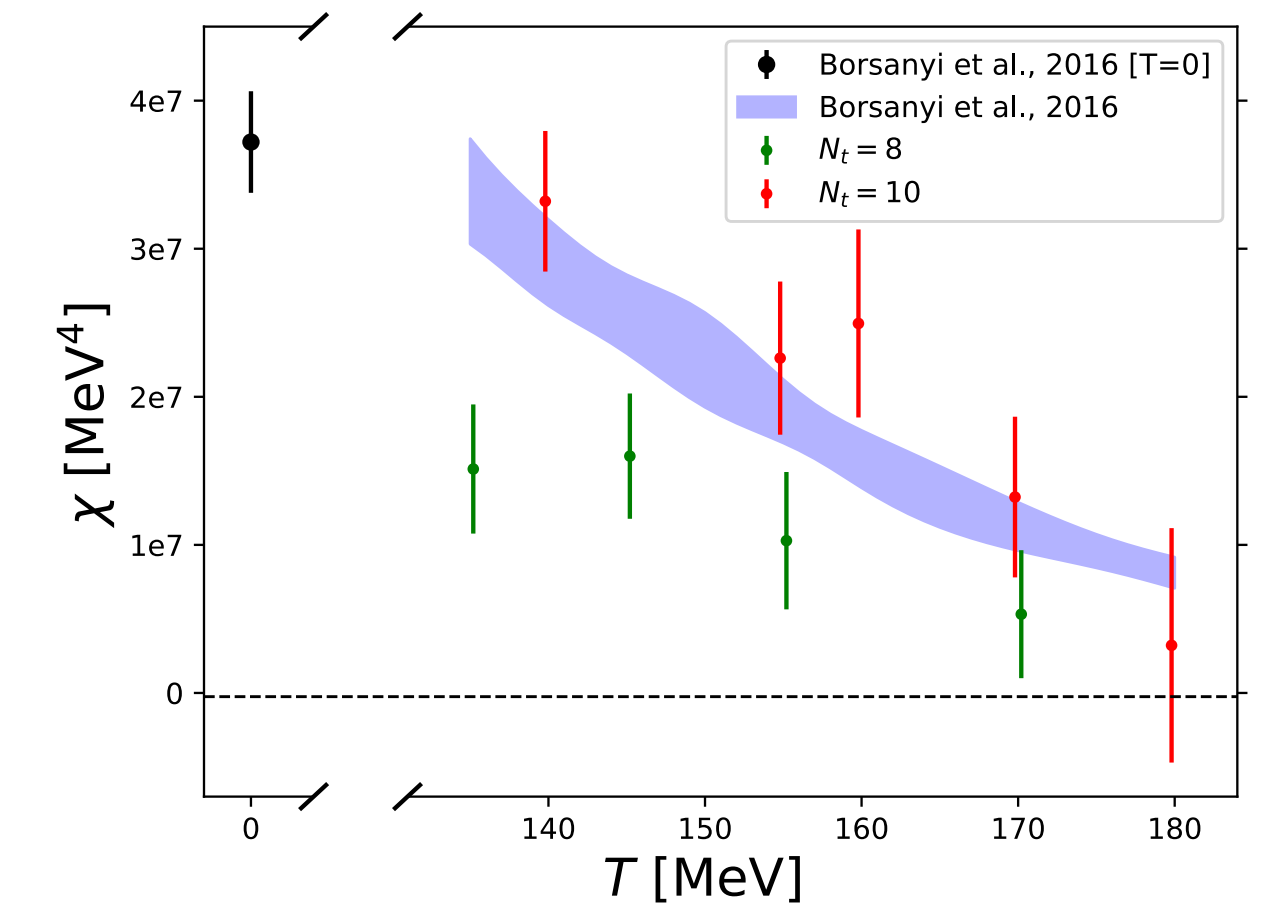
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- Dirac spectrum: **peak at  $\rho(\lambda \rightarrow 0)$**   
for  $N_s/N_t \gtrsim 4 - 5$  at  $T \gtrsim T_{pc}$
- Near-zero modes become **localised** around  $T_c$



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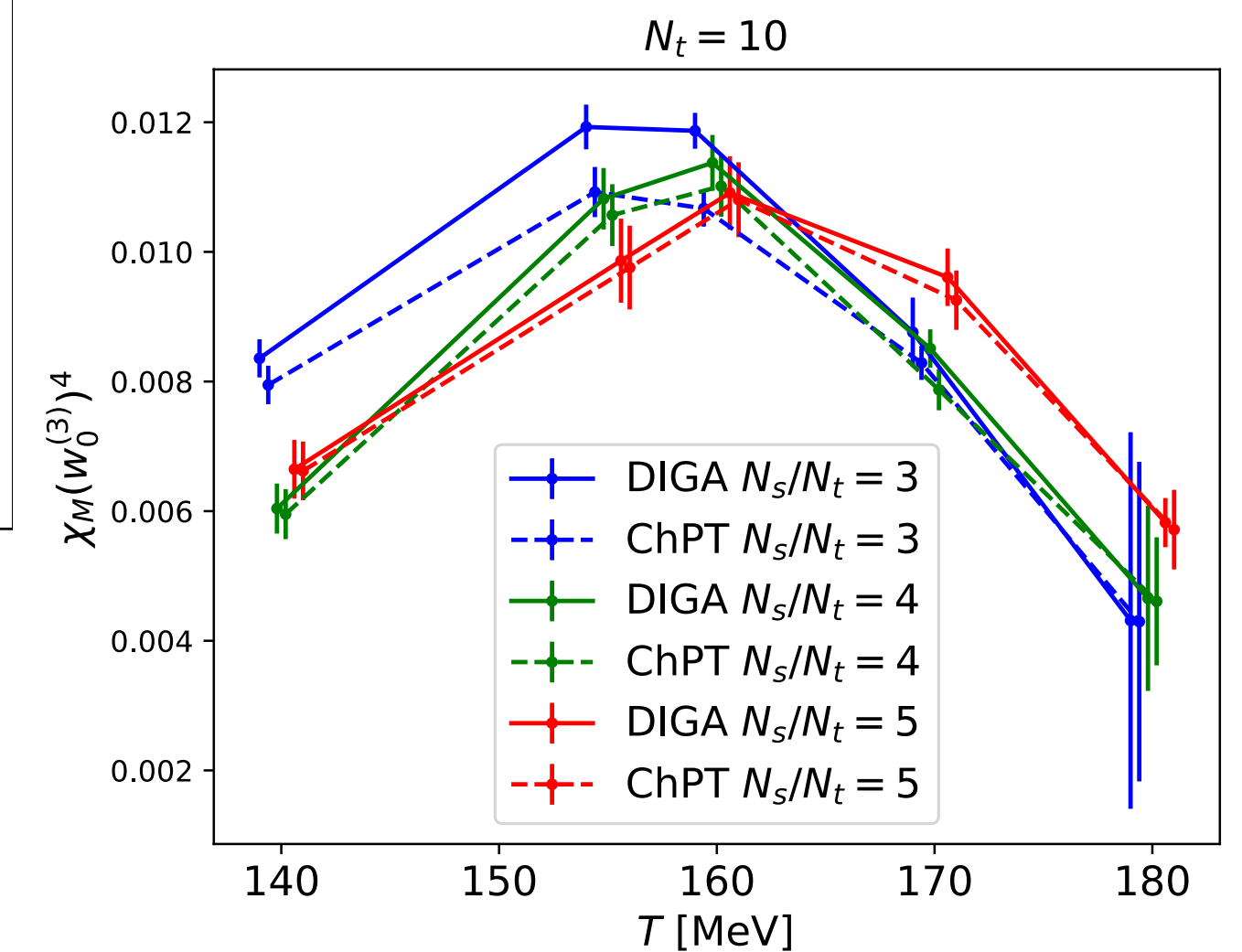
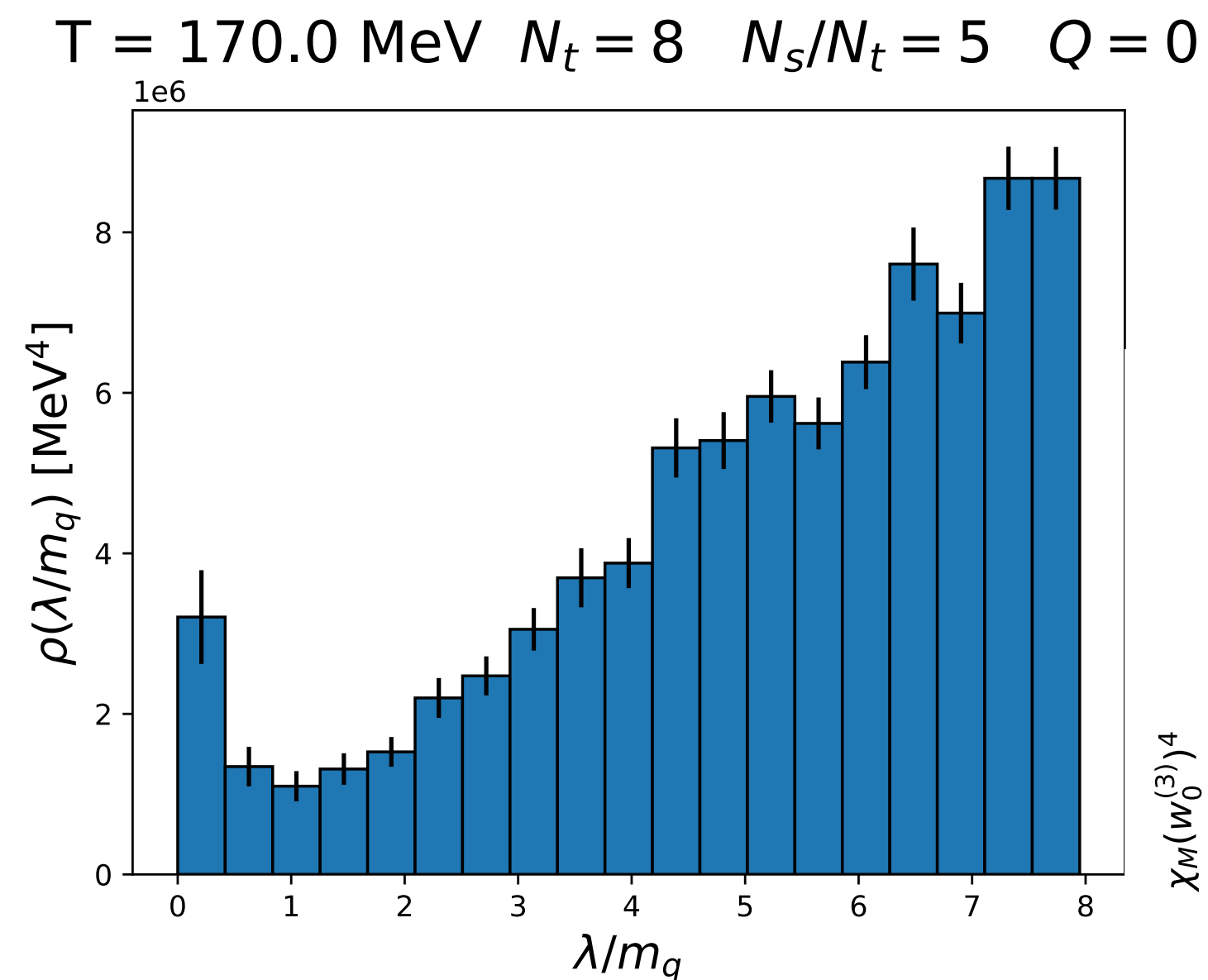
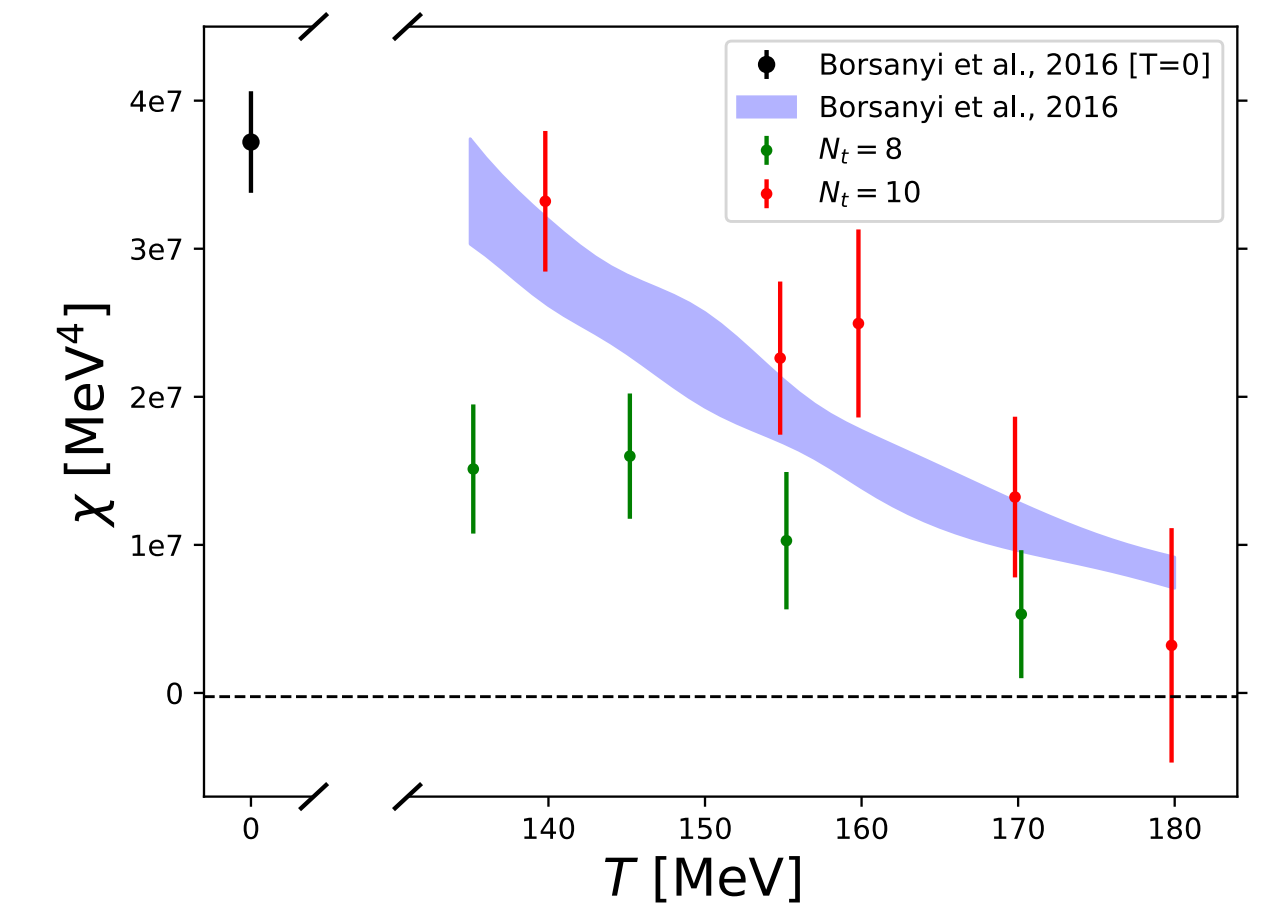
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for  $N_s/N_t \gtrsim 4 - 5$  at  $T \gtrsim T_{pc}$
- Near-zero modes become **localised** around  $T_c$
- Cutoff effects for  $N_t = 8$ , larger  $N_t$  ?



# Summary

Purely overlap result!

- **Dynamical overlap fermions** at  $m_\pi = m_\pi^{\text{phys}}$ , summed over  $Q$
- $\chi_Q$  from overlap simulations
- Chiral observables around  $T_c$
- Consistent with chiral crossover
- Dirac spectrum: **peak at  $\rho(\lambda \rightarrow 0)$**   
for  $N_s/N_t \gtrsim 4 - 5$  at  $T \gtrsim T_{pc}$
- Near-zero modes become **localised** around  $T_c$
- Cutoff effects for  $N_t = 8$ , larger  $N_t$  ?



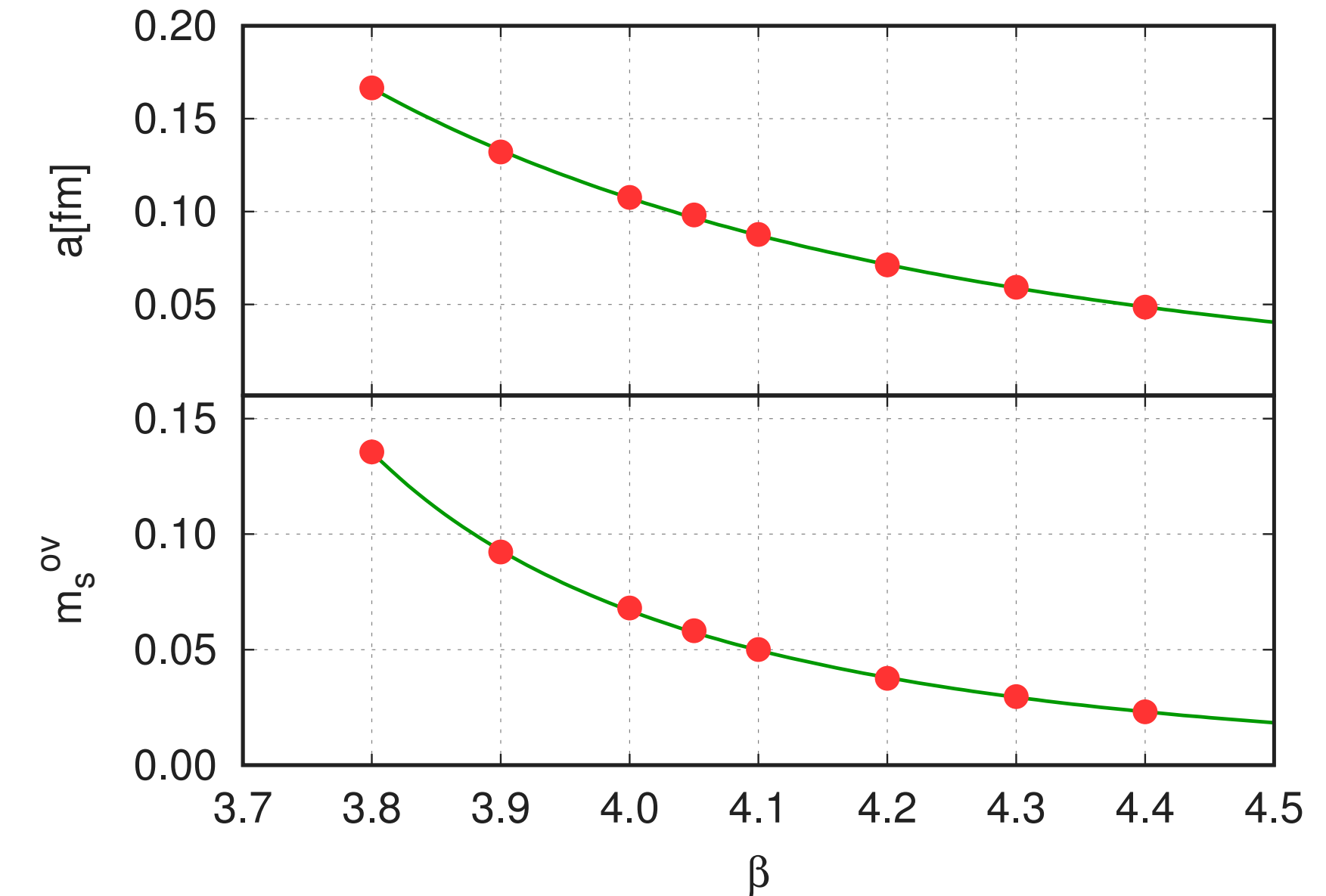
Thank you for your attention!

**Backup**

# Lattice details, scale setting

## Scale setting from simulations with large $m_\pi$

- Simulations are done along the LCP
- Scale setting: require  $T = 0$  simulations
- $N_f = 3$  staggered simulations,  $T = 0$ ,  $w_0^{(3)} = 0.153(1)$  fm,  $m_\pi^{(3)} = 712(5)$  MeV
- $N_f = 3$  overlap simulations,  $T = 0$ , at each  $\beta$  tune  $m_s^{\text{ov}}$  to have  $m_\pi w_0 \equiv m_\pi^{(3)} w_0^{(3)}$
- $N_f = 2 + 1$  overlap simulations,  $T \neq 0$ :  $m_s = m_s^{\text{ov}}$ ,  $m_{ud} = R m_s^{\text{ov}}$ ,  $a = w_0^{(3)} / w_0^{\text{ov}}$
- Physical point:  $m_{ud} = m_{ud}^{(\text{phys})}$ ,  $m_s = m_s^{(\text{phys})}$



[Borsanyi et al., 2016]

# Implementing odd number of flavours

Exploiting  $Q = \text{const}$

- Monte Carlo: determinant of a **hermitian** operator  $H^2 = D_{\text{ov}} D_{\text{ov}}^\dagger$ :  $N_f = 2$
- To simulate  $N_f = 1$  (strange quark): need to take the **square root**
- Chirality projectors:  $P_\pm = \frac{1 \pm \gamma_5}{2}$ ,  $H_\pm^2 = P_\pm H^2 P_\pm$
- Fixed topology  $Q = \text{const}$ :  
 $\det H^2 \sim \det H_+^2 \det H_-^2 \sim (\det H_+^2)^2 \sim (\det H_-^2)^2$
- Take  $\det H_+^2$  or  $\det H_-^2$